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UPS Photoemission Properties of (SN)_x. P. MENGEL,
W. D. GROBMAN, I. B. ORTENBURGER, P. M. GRANT and B. H.
SCHECHTMAN, IBM Research Laboratory, San Jose, Calif.--
We have performed UPS measurements on in situ deposited
films of (SN)_x as a function of incident photon energy
in the range 7.6 - 40.8 eV. The results are in general
agreement with previously reported XPS studies.¹ Re-
laxation effects lower the observed photoemissive yield
at the Fermi level to zero. We analyze our data through
comparison with theoretical energy distribution curves
obtained from pseudopotential eigenvalues and eigen-
functions. We also report some initial results on the
UPS spectrum of partially polymerized S₂N₂.

¹P. Mengel, P. M. Grant, W. E. Rudge, B. H. Schechtman
and D. W. Rice, Phys. Rev. Letters 35, 1803 (1975).

Submitted by

Signature of APS Member

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BE 4 Transport Properties of $Hg_{2.86}(AsF_6)_2$. R. Spal, T. Wei, C.K. Chiang, A.F. Garito and A.J. Heeger, N. Miro and A.G. MacDiarmid, University of Pennsylvania, Phila., Penna.--We report the temperature dependent dc electrical conductivity of the novel conductor $Hg_{2.86}(AsF_6)_2$ which consists of planes of interpenetrating one dimensional chains of Hg atoms. The measured conductivity parallel to the planes is $3.5 \times 10^3 (\Omega\text{-cm})^{-1}$ at room temperature and exhibits a negative temperature coefficient down to 4.2K with an anomaly near 220K. Anisotropy measurements using the Montgomery method will be reported in addition. The conductivity measurements will be compared with studies of a plasma edge in the visible optical reflectance and measurements using DSC.

BE 5 Anionic Linear Chain Iridium Carbonyl Halides. A. P. GINSBERG, J. W. KOEFLKE, J. J. HAUSER, K. W. WEST, F. J. DI SALVO, C. R. SPRINKLE and R. L. COHEN, Bell Laboratories.--The compounds $K_{0.60}Ir(CO)_2Cl_2 \cdot 0.5H_2O$, (TTF) $_{0.61}Ir(CO)_2Cl_2$ (TTF = Tetraphthalvalenium), $K_{0.57}Ir(CO)_2Br_2 \cdot 0.2CH_3COCH_3$ and $Cs_{0.60}Ir(CO)_2Br_2$ have been investigated by chemical analysis, infrared and ^{193}Ir Mössbauer spectroscopy, electrical conductivity and magnetic susceptibility measurements. Conducting linear chains of $cis-[Ir(CO)_2X_2]^{+0.00}$ ($X = Cl, Br$) units are shown to be present in these compounds. The apparently well defined compounds $K_{0.68}Ir(CO)_2Cl_2 \cdot 4.2 \cdot 0.2CH_3COCH_3$ and $Na_{0.93}Ir(CO)_2Cl_2 \cdot 32 \cdot 0.3CH_3COCH_3$, are reformulated as $K_{0.60}Ir(CO)_2Cl_2 \cdot 0.14KCl \cdot 0.2CH_3COCH_3$ and $Na_{0.61}Ir(CO)_2Cl_2 \cdot 0.32NaCl \cdot 0.3CH_3COCH_3$, in which the KCl and NaCl are present interstitially in the lattice of linear chains.

BE 6 EPR pair spectra in a dimeric chromium phosphinate.* P. D. KRASICKY, A. L. RITTER, J. C. SCOTT and R. H. SILSBEE, Cornell University--The dimeric material, di- μ -diphenylphosphinatoacetylacetonatochromium (III)⁽¹⁾, consists of molecules in which the two chromium atoms, each of spin $S_{1,2} = 3/2$ are linked by phosphinate ($-OP_2O-$) bridges, and is therefore the simplest member of the family of one-dimensional poly(chromium phosphinates).⁽²⁾ The "endcapping" acetylacetone ligands complete the closely octahedral environment of six oxygen atoms around each chromium. Within the molecule the chromium-chromium separation is 5 Å, compared to ~ 10 Å between molecules. The EPR pair spectrum at X band indicates an antiferromagnetic exchange, $J \approx 4$ K, comparable to that in the polymers. At temperatures $T \leq 4$ K the S=3 (total spin) multiplet is virtually unpopulated, and we have identified transitions within the S=1 and S=2 multiplets. The data are interpreted in terms of a spin Hamiltonian derived by coupling the single ion Hamiltonians with an exchange term.

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1C.E. Wilkes and R.A. Jacobson, Inorg. Chem. 4, 99 (1965).
2J.C. Scott, T.S. Wei, A.F. Garito, A.J. Heeger, H.D. Gillman and P. Nannelli, Proc. Conf. on Magnetism and Magnetic Materials, Phila., 1975 (to be published).

BE 7 UPS Photoemission Properties of $(SN)_x$. P. MENGEI, W. D. GROBMAN, I. B. ORTENBURGER, P. M. GRANT and B. H. SCHECHTMAN, IBM Research Laboratory, San Jose, Calif.--We have performed UPS measurements on in situ deposited films of $(SN)_x$ as a function of incident photon energy in the range 7.6 - 40.8 eV. The results are in general agreement with previously reported XPS studies. Relaxation effects lower the observed photoemissive yield at the Fermi level to zero. We analyze our data through comparison with theoretical energy distribution curves obtained from pseudopotential eigenvalues and eigenfunctions. We also report some initial results on the UPS spectrum of partially polymerized S_2N_2 .

1P. Mengel, P. M. Grant, W. E. Rudge, B. H. Schechtman and D. W. Rice, Phys. Rev. Letters 35, 1803 (1975).

BE 8 XPS Core-Level Spectroscopy, Charge Transfer and Electrostatic Interactions in $(SN)_x$ and Other SN Allotropes. A.J. EPSTEIN, W.R. SALANECK, N.O. LIPARI and J.W.-p LIN, Xerox Webster Research Center, Webster, N.Y. 14580.--The binding energies of S_{2p} and N_{1s} levels in $(SN)_x$ have been studied via x-ray photoemission spectroscopy and results compared with neutral S_8 and N_2 molecules respectively, as well as S_4N_4 . The effective charge transfer from sulfur to nitrogen is $0.5|e|$ in $(SN)_x$ as compared to $0.6|e|$ in S_4N_4 . Shake-up structure on the core level lines of $(SN)_x$ is assigned to excitation of conduction electron plasmons. Using the above XPS measured net charges, the electrostatic interactions in S_2N_2 , S_4N_4 and $(SN)_x$ were calculated. Though the total electrostatic potential at any site due to all other atoms in a crystal is large and nearly identical for all three allotropes, the net electrostatic interaction between molecules (polymer chains) is small and nearly identical in all three cases.

BE 9 Thermorelectance of Organic and Polymeric Metals. P. M. GRANT, P. MENGEI, E. M. ENGLER, G. CASTRO and G. B. STREET, IBM Research Laboratory, San Jose, Calif.--We have examined the thermorelectance spectrum of $(SN)_x$ crystals and films and of $(TTF)(TCNQ)$ and $(TSeF)(TCNQ)$ crystals in energy ranges appropriate to their Drude properties and low-lying interband transitions as a function of temperature from 300°K to 15°K. In $(SN)_x$, we observe a red shift of the Drude edge which we associate with movement of the Fermi level due to volume effects. Perov and Fischer¹ have reported polarization effects in the thermorelectance spectrum of $(TTF)(TCNQ)$ connected with the Peierls-Fröhlich transition. We have made similar measurements on $(TSeF)(TCNQ)$ and compare our results with theirs for $(TTF)(TCNQ)$.

1P. I. Perov and J. E. Fischer, Phys. Rev. Letters 33, 521 (1974).

BE 10 Calculated Optical Properties of $(SN)_x$. I. B. ORTENBURGER, W. E. RUDGE and P. M. GRANT, IBM Research Laboratory, San Jose, Calif.--The tensorial dielectric constant $\epsilon_2(\omega)$ and the tensorial photoemission yield $D(E, \omega)$ have been calculated for $(SN)_x$ using a pseudopotential interpolation scheme based^x on a first principles OPW band structure. $D(E, \omega)$ and $\epsilon_2(\omega)$ were computed by the Gilat-Raubenheimer zonal integration method. Our results will be compared to experimental photoemission and optical¹ measurements on randomly oriented $(SN)_x$ films.

1A. A. Bright and A. F. Garito, to be published.

BE 11 Molecular Cluster and Band Structure Calculations for $(SN)_x$. D. R. SALAHUB and R. P. MESSMER, General Electric.--Molecular cluster and band structure calculations for $(SN)_x$ show that (i) a single $(SN)_x$ chain should be a Peierls insulator, (ii) semi-metallic behavior is brought about by interchain coupling which moves the Fermi level from the Peierls gap to a place where it intersects two bands, (iii) the most important interactions occur in a $\bar{1}00$ plane, and not as previously assumed in a $\bar{1}02$ plane. Important aspects of the electronic structure of intermediates involved in the polymerization and sublimination of $(SN)_x$ are also discussed.

BE 12 Band Structure, Density of States, and Fermi Surface Topology of the $(SN)_x$ Crystal.* W. Y. CHING, J. G. HARRISON, and CHUN C. LIN, U. of Wisconsin

Photoemission talk

We carried out x-ray and UPS photoemission measurements on $(SN)_x$ films to reveal experimentally the electronic properties of this compound and compare it with theoretical calculations.

In this talk we would like to concentrate on the result for the valenceband region of $(SN)_x$.

The first slide gives an overall picture of the valenceband region of $(SN)_x$. The upper part shows the experimental data from XPS measurement (Al $K\alpha$ 1.4 keV line, dots raw data, solid line inelastic background subtracted). The lower part gives a comparison with 2 different theoretical DOS, IDTB ($T \& P$) and 3DOPW ($R \& G$). The valenceband extends about 30 eV below E_F , and we find a good qualitative agreement between both theoretical calculation and the experimental data, and in addition the 3DOPW gives a good quantitative agreement in peak energy position and relative peak intensity too.

In the UPS measurements we hoped with better resolution to clarify the structure near the Fermi level, density of state appears to be low, the exactly energy position of peak 1, which shows the coincidence in energy with the DOPW only for the Boudelle crystal structure.

The next slide shows EDC's obtained from a continuously H₂ lamp for eleven photons between 7.6 and 11.1 eV. The cross line indicates zero kinetic energy. The solid curves are the experimental results, dashed are theoretical EDC calculations.

The films have been prepared in-situ of the UHV chamber. After the measurement we took them out for X-ray analysis, which shows the (102) plane in a rotating-crystal method type of experiment. This gives evidence for a preferred orientations in the films but

there may be still a disorder in the direction perpendicular to the chain axes.

The EDC's are in general agreement with XPS measurement. The peaks at 1.0 eV and 4.5 eV correspond to peak 1 and 2 in the slide before, relative intensity of the second peak appreciable higher, and there is an additional structure at about 2.5 eV. Surprising again is the low density of state at the Fermi level (nearly zero).

The theoretical EDC's (dashed line) are calculated from first principle pseudopotential eigenvalues and eigenfunctions. The experimentally found work function of 4.2 eV was used for applying escape probability function on each of the theoretical EDC's.

The theoretical calculations are in reasonably agreement with the experimental data.

the lower photon energies. All structure in the exp. curves present in the calculation, emphasize that we used the Penn structure, which we believe is the more appropriate structure, but the correlation in energy position of peak 1 is much better with the Baudelle structure. Most important is that also the calculate EDC show a very low density of state at E_F . Now we can clearly show that this is due to Matrix element effect, which gives a very low transition probability near the Fermi Level, about $\frac{1}{5}$ of the average matrix element, at higher photon energies. The increase of the DOS¹ in the theoretical calculation is not found in the experimental data. In the theory the reason is artificial final state effect and we believe that the pseudopotential method is ^{just} insufficient in determining the conduction bands far above

the Fermi level.

Indeed measurements with higher photon-energies, which are shown in the next slide for 10.2, 16.8, 21 and 40.8 eV show again a low density of states at the Fermi level. Again we see the same structure obtained at the other photon energies, and the bottom curves give a comparison with a pseudopotential DOS.

Instead of a increase of the DOS at E_F , we even get a gap of 0.2 eV at the Fermi level; and there are some reason to believe, that the film exist in a non fully developed crystalline phase - presumably amorphous considering the order perpendicular to the chain axes. Reasons: conductivity, reflection measurement, structure analysis.

In this case the films are most disordered at the surface and therefore depending on the sample parallel to the

Photoemission experiment, as gap opens
at photon energies where the escape length
has it's minimum.





