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One-Dimensional Conductors

Thermoreflectance of Organic and Polymeric Metals. P. M. GRANT, P. MENGEL, E. M. ENGLER, G. CASTRO and G. B. STREET, IBM Research Laboratory, San Jose, Calif.—We have examined the thermoreflectance spectrum of (SN) crystals and films and of (TTF)(TCNQ) and (TSeF)(TCNQ) crystals in energy ranges appropriate to their Drude properties and low-lying interband transitions as a function of temperature from 300°K to 15°K. In (SN), we observe a red shift of the Drude edge which we associate with movement of the Fermi level due to volume effects. Perov and Fischer have reported polarization effects in the thermoreflectance spectrum of (TTF)(TCNQ) connected with the Peierls-Fröhlich transition. We have made similar measurements on (TSeF)TCNQ) and compare our results with theirs for (TTF)(TCNQ).

¹P. I. Perov and J. E. Fischer, Phys. Rev. Letters <u>33</u>, 521 (1974).

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BE 4 Transport Properties of Hg2.86(AsF)6, R. Spal, T. Wei, C.K. Chiang, A.F. Garito and A.J. Heeger, N. Miro and A.G. MacDiarmid, University of Pennsylvania, Phila., Penna. --We report the temperature dependent dc electrical conductivity of the novel conductor Hg2.86(AsF6)2 which consists of planes of interpenetrating one dimensional chains of Hg atoms. The measured conductivity parallel to the planes is 3.5 x $10^3 (\Omega - \mathrm{cm})^{-1}$ at room temperature and exhibits a negative temperature coefficient down to 4.2K with an anomaly near 220K. Anisotropy measurements using the Montgomery method will be reported in addition. The conductivity measurements will be compared with studies of a plasma edge in the visible optical reflectance and measurements using DSC.

BE 5 Anionic Linear Chain Iridium Carbonyl Halides. A. P. GINSBERG, J. W. KOEPKE, J. J. HAUSER, K. W. WEST, F. J. DI SALVO, C. R. SPRINKLE and R. L. COHEN, Bell Laboratories. -- The compounds Ko.60Ir(CO)2Cl2.0.5H2O, (TTF)0.61Ir(CO)2Cl2 (TTF = Tetrathiafulvalenium), Ko.57Ir(CO)2Br2.0.2CH3COCH3 and Cso.60Ir(CO)2Br2 have been investigated by chemical analysis, infrared and 193Ir Mossbauer spectroscopy, electrical conductivity and magnetic susceptibility measurements. Conducting linear chains of cis-[Ir(CO)2X2]-0.50 (X = Cl, Br) units are shown to be present in these compounds. The apparently well defined compounds Ko.98Ir(CO)2Cl2.42.0.2CH3COCH3 and Nao.93Ir(CO)2Cl2.32.0.3CH3COCH3, are reformulated as Ko.60Ir(CO)2Cl2.32.0.3CH3COCH3, in which the KCl and NaCl are present interstitially in the lattice of linear chains.

BE 6 EPR pair spectra in a dimeric chromium phosphinate.* P. D. KRASICKY, A. L. RITTER, J. C. SCOTT and R. H. SILSBEE, Cornell University -- The dimeric matedi-µ-diphenylphosphinatoacetylacetonatochromium (III)(1), consists of molecules in which the two chromium atoms, each of spin $S_{1,2}=3/2$ are linked by phosphinate (-OP Φ_2 O-) bridges, and is therefore the simplest member of the family of one-dimensional poly(chromium phosphinates).(2) The "endcapping" acetylacetonate ligands complete the closely octahedral environment of six oxygen atoms around each chromium. Within the molecule the chromium-chromium separation is 5Å, compared to ~10Å between molecules. The EPR pair spectrum at X band indicates an antiferromagnetic exchange, J 4K, comparable to that in the polymers. At temperatures T4K the S=3 (total spin) multiplet is virtually unpopulated, and we have identified transitions within the S=1 and S=2 multiplets. The data are interpreted in terms of a spin Hamiltonian derived by coupling the single ion Hamiltonians with an exchange term.

*Supported by the National Science Foundation.

1C.E. Wilkes and R.A. Jacobson, Inorg. Chem.4,99(1965).

2J.C. Scott, T.S. Wei, A.F. Garito, A.J. Heeger, H.D. Gillman and P. Nannelli, Proc. Conf. on Magnetism and Magnetic Materials, Phila., 1975 (to be published)

BE 7 UPS Photoemission Properties of (SN). P. MENGEL, W. D. GROBMAN, I. B. ORTENBURGER, P. M. GRANT and B. H. SCHECHIMAN, IBM Research Laboratory, San Jose, Calif.—We have performed UPS measurements on in situ deposited films of (SN) as a function of incident photon energy in the range 7.6 - 40.8 eV. The results are in general agreement with previously reported XPS studies. Relaxation effects lower the observed photoemissive yield at the Fermi level to zero. We analyze our data through comparison with theoretical energy distribution curves obtained from pseudopotential eigenvalues and eigenfunctions. We also report some initial results on the UPS spectrum of partially polymerized S₂N₂.

P. Mengel, P. M. Grant, W. E. Rudge, B. H. Schechtman and D. W. Rice, Phys. Rev. Letters 35, 1803 (1975). BE 8 XPS Core-Level Spectroscopy, Charge Transfer and Electrostatic Interactions in (SN) $_{\rm X}$ and Other SN Allotropes. A.J. EPSTEIN, W.R. SALANECK, N.O. LIPARI and J.W-p LIN, Xerox Webster Research Center, Webster, N.Y. 14580.—The binding energies of S2p and N1 $_{\rm S}$ levels in (SN) $_{\rm X}$ have been studied via X-ray photoemission spectroscopy and results compared with neutral S8 and N2-molecules respectively, as well as S4N4. The effective charge transfer from sulfur to nitrogen is 0.5|e| in (SN) $_{\rm X}$ as compared to 0.6|e| in S4N4. Shake-up structure on the core level lines of (SN) $_{\rm X}$ is assigned to excitation of conduction electron plasmons. Using the above XPS measured net charges, the electrostatic interactions in S2N2, S4N4 and (SN) $_{\rm X}$ were calculated. Though the total electrostatic potential at any site due to all other atoms in a crystal is large and nearly identical for all three allotropes, the net electrostatic interaction between molecules (polymer chains) is small and nearly identical in all three cases.

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BE 9 Thermoreflectance of Organic and Polymeric Metals. P. M. GRANT, P. MENGEL, E. M. ENGLER, G. CASTRO and G. B. STREET, IBM Research Laboratory, San Jose, Calif.—We have examined the thermoreflectance spectrum of (SN) crystals and films and of (TTF)(TCNQ) and (TSeF)(TCNQ) crystals in energy ranges appropriate to their Drude properties and low-lying interband transitions as a function of temperature from 300°K to 15°K. In (SN), we observe a red shift of the Drude edge which we associate with movement of the Fermi level due to volume effects. Perov and Fischer have reported polarization effects in the thermoreflectance spectrum of (TTF)(TCNQ) connected with the Peierls-Fröhlich transition. We have made similar measurements on (TScF)TCNQ) and compare our results with theirs for (TTF)(TCNQ).

¹P. I. Perov and J. E. Fischer, Phys. Rev. Letters <u>33</u>, 521 (1974).

BE 10 Calculated Optical Properties of (SN). I. B. ORTENBURGER, W. E. RUDGE and P. M. GRANT, IBM Research Laboratory, San Jose, Calif.—The tensorial dielectric constant $\epsilon_2(\omega)$ and the tensorial photoemission yield $D(E,\omega)$ have been calculated for (SN) using a pseudopotential interpolation scheme based on a first principles OPW band structure. $D(E,\omega)$ and $\epsilon_2(\omega)$ were computed by the Gilat-Raubenheimer zonal integration method. Our results will be compared to experimental photoemission and optical measurements on randomly oriented (SN) films.

A. A. Bright and A. F. Garito, to be published.

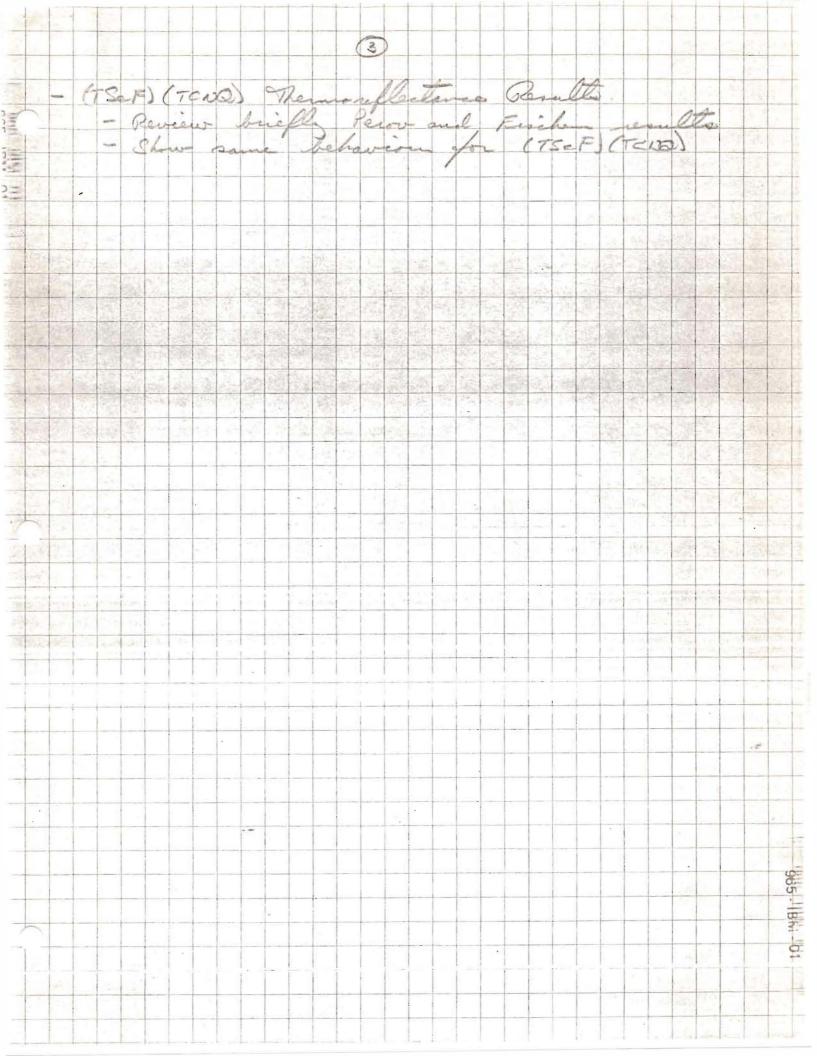
BE11 Molecular Cluster and Band Structure Calculations for (SN) . D. R. SALAHUB and R. P. MESSMER, General Electric: --Molecular cluster and band structure calculations for (SN) x show that (i) a single (SN) x chain should be a Peierls insulator, (ii) semi-metallic behavior is brought about by interchain coupling which moves the Fermi level from the Peierls gap to a place where it intersects two bands, (iii) the most important interactions occur in a 100 plane, and not as previously assumed in a 102 plane. Important aspects of the electronic structure of intermediates involved in the polymerization and subliniation of (SN) x are also discussed.

BE 12 Band Structure, Density of States, and Fermi Surface Topology of the $(SN)_x$ Crystal.* W. Y. CHING, J. G. HARRISON, and CHUN C. LIN, U. of Wisconsin-

Outline of BE9: Thermoreflectance of Organic and Polymer metalo" Thermore flectance is a form of derivative applied spectroscopy which has frequently been applied to conventional metals to help elicidate details of their electronic structure difficult to consiste of indirectly heating the sample via detecting the impressed modulation in the In this talk we will discuss mostly (SN) x and if time remute, some of our results for (TSEF) (TCNQ) as well. Our (SU), tudies were motivated by our finding a very weak temperature dependent of (SN), cuptales. 13 - Change in Ode of 6 observed over the same temorature range. - Change of Tode of 66 116 for P > 10 kbas opical grozenties Equations for themoreflestance For money metals, suplat dominates although reason to assume some for (SN). no a puri reflectivity with Temperature will setermine the sign of DR/R.

(SN) Film Thermorellectores in Duide Rouge Show red shift in both sign and derection of plasma query, connection with Tide now Reasons why (SN) oftical projection do not - Ensue probably lies in the difference in the two lifetime soverning de and optical. Relatively strong optical shoron bands lie in the wavelength range 5-10 u. It is reasonable to come stong temperature independent dampin from these glemon could be the doneinant factor at optical presences.

- If the de centlerin is doneinated by electron - electrons effects as an expected by the chians and coworkers, their interaction could be greath affected by changes in the shape of Farmi surface lue to dilation, particula in the interchain disctions, which could also graduce the small optical red shift in change energy yet lawe the ortical lifetime affected. He will be most satisfaction of but Effects in (SN) x Film Thermore lection - De not arise from E 116 to structur De ain from E/b (700) er structure - Week tomperature dependence of we and Tomore suggested Toc Temperature Sepandence completely doninated by different no.



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THERMOREFLECTANCE EQUATIONS

$$\frac{1}{2} \frac{\Delta R}{\Delta T} = \left(\frac{\partial \Omega_1 R}{\partial \epsilon_1}\right) \frac{\Delta \epsilon_2}{\Delta T} + \left(\frac{\partial \Omega_1 R}{\partial \epsilon_2}\right) \frac{\Delta \epsilon_2}{\Delta T}$$

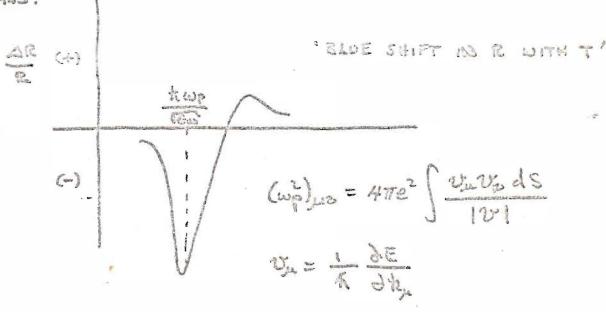
$$\beta_1 (\epsilon_1, \epsilon_2) \qquad \beta_2 (\epsilon_1, \epsilon_2)$$

$$\frac{\Delta \mathcal{E}_{1,2}}{\Delta T} = \frac{\partial \mathcal{E}_{1,2}}{\partial \omega_0} \frac{\Delta \omega_0}{\Delta T} + \frac{\partial \mathcal{E}_{1,2}}{\partial \omega_0} \frac{\Delta \omega_0}{\Delta T}$$

$$+ \frac{\partial \mathcal{E}_{1,2}}{\partial \omega_0} \frac{\Delta T}{\Delta T} + \frac{\partial \mathcal{E}_{1,2}}{\partial \mathcal{E}_{00}} \frac{\Delta \omega_0}{\Delta T}$$

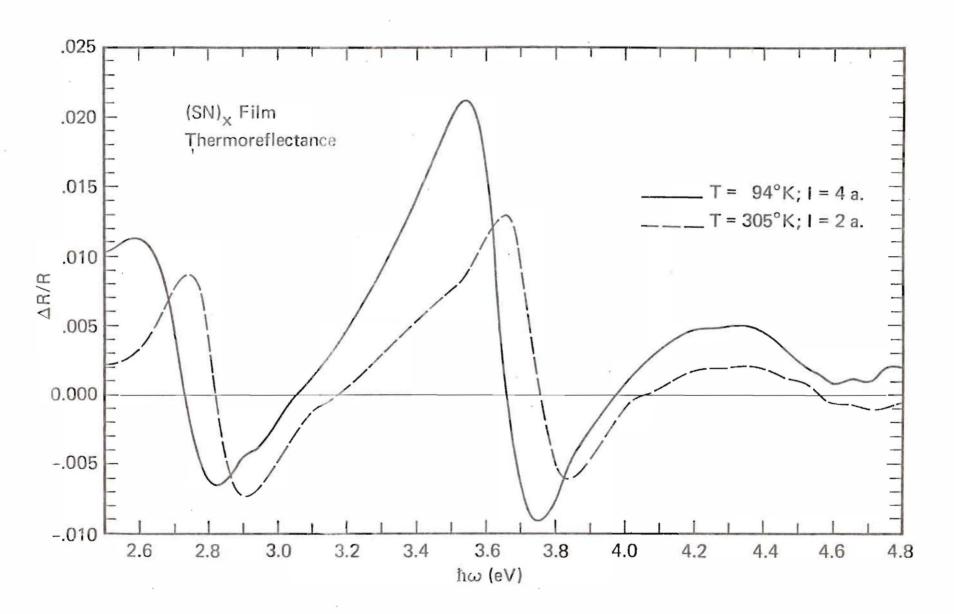
$$E(\omega) = E\omega + \omega^2 - \omega^2 + \omega / v$$

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