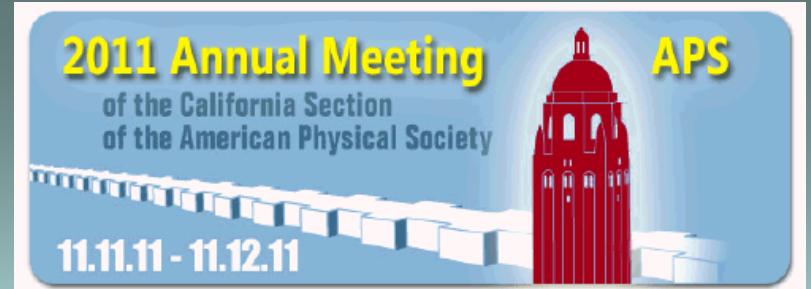




California APS

2011 Annual Meeting



SLAC NATIONAL ACCELERATOR LABORATORY

The Great Quantum Conundrum

Paul M. Grant
W2AGZ Technologies
www.w2agz.com

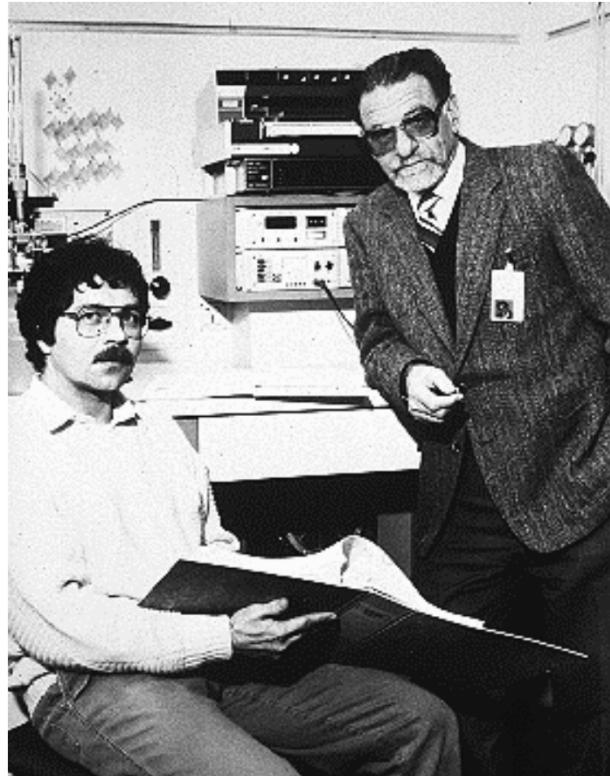
AGING IBM PENSIONER

SLAC
2575 Sand Hill Road
Menlo Park, CA 94025 USA

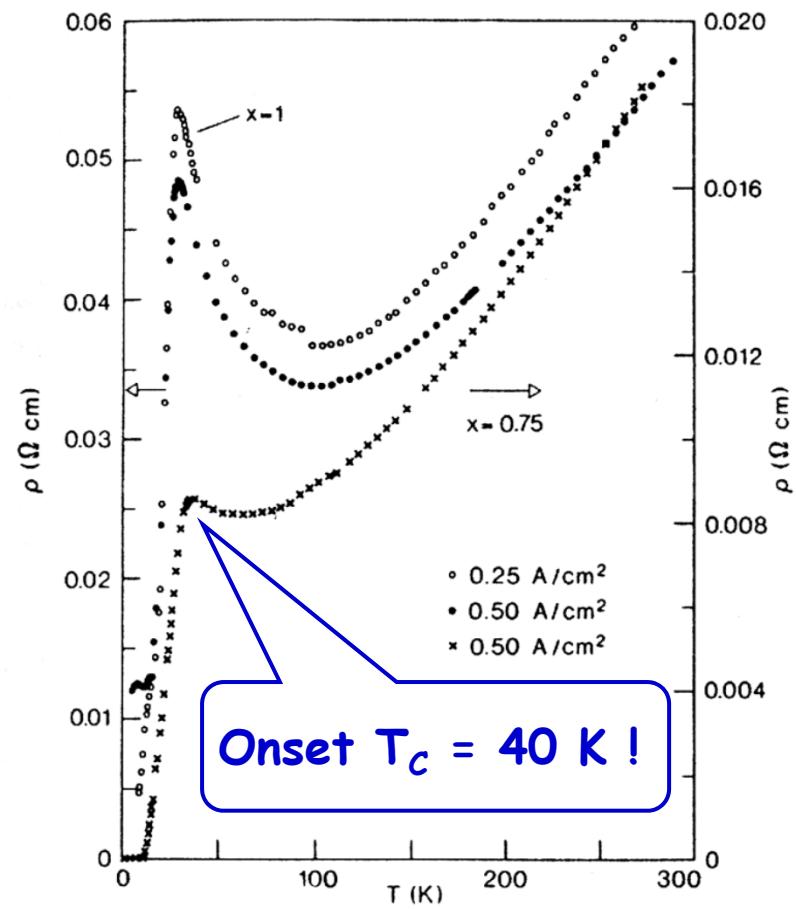
SLAC
Sand Hill Road
Menlo Park, CA

1986

Another Big Surprise!



Bednorz and Mueller
IBM Zuerich, 1986



1987

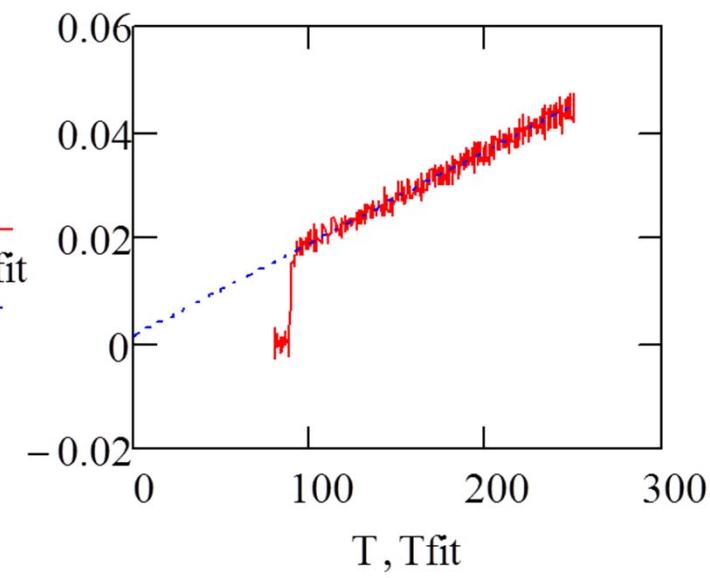
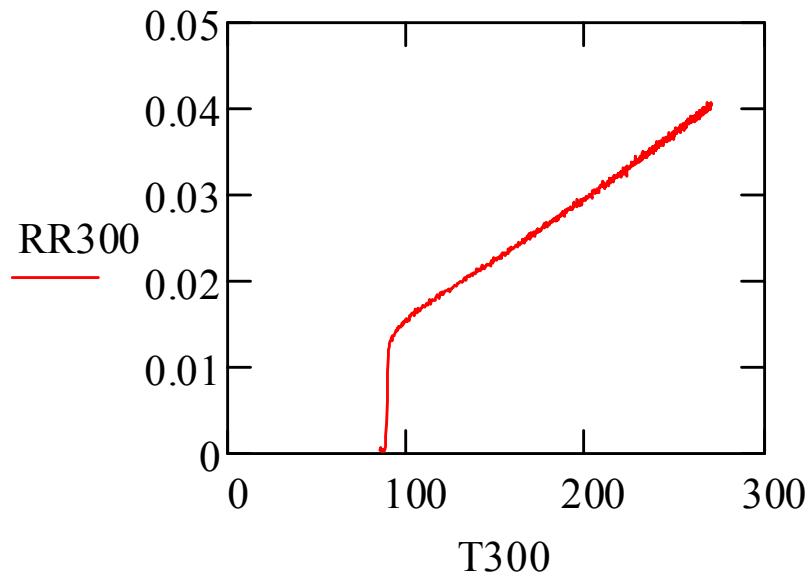
“The Prize!”



Associated Press

J. Georg Bednorz, left, and K. Alex Müller after learning they had won the Nobel Prize in physics.

2 Get Nobel for Unlocking Superconductor Secret



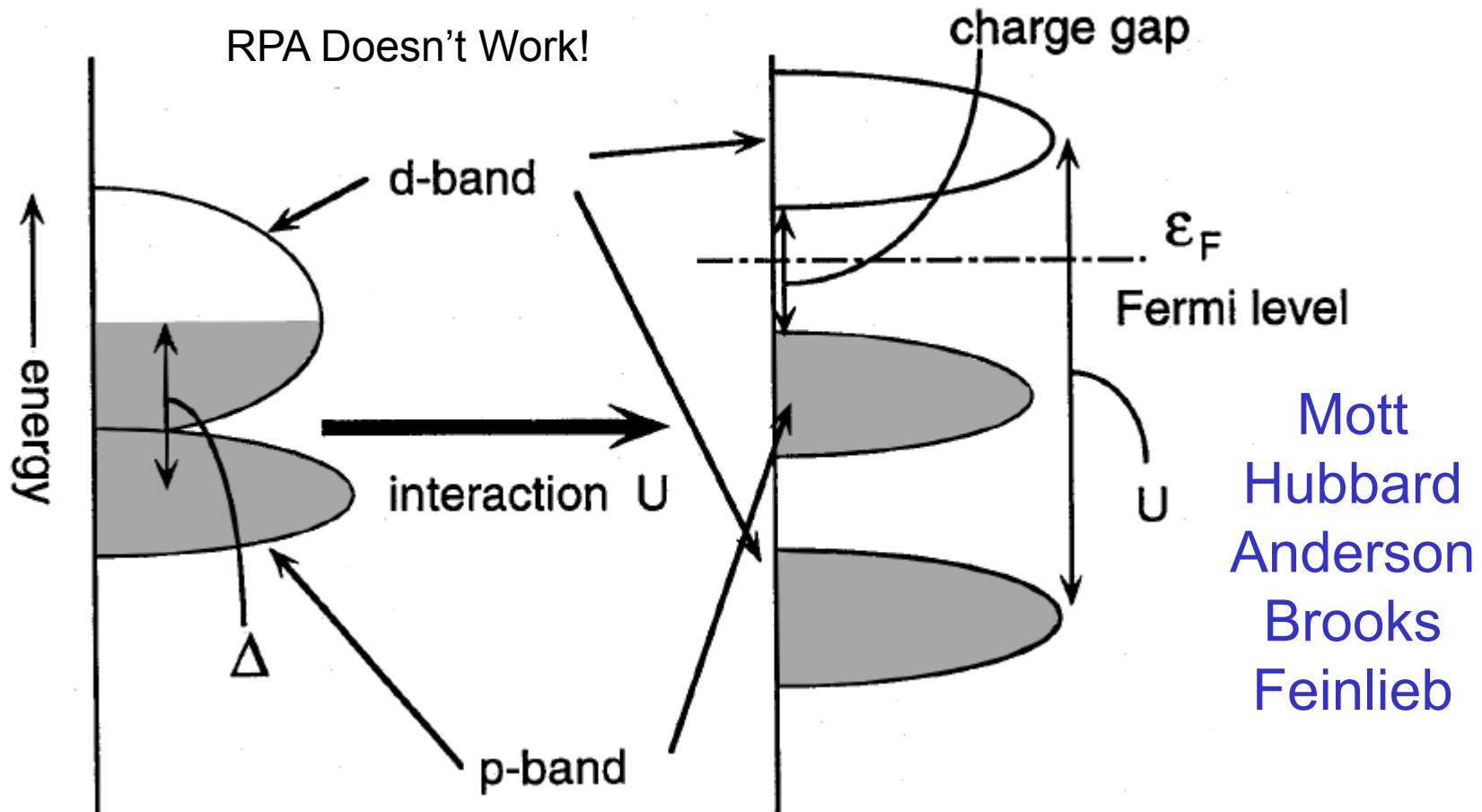
A quarter-century after Bednorz-Mueller, it is still not clear why:

- 1) T_c is so high
 - 2) The normal state appears non-Fermi liquid-like
- That is the Great Quantum Conundrum

Transition Metal Oxides

“Should be Metals, But Aren’t”

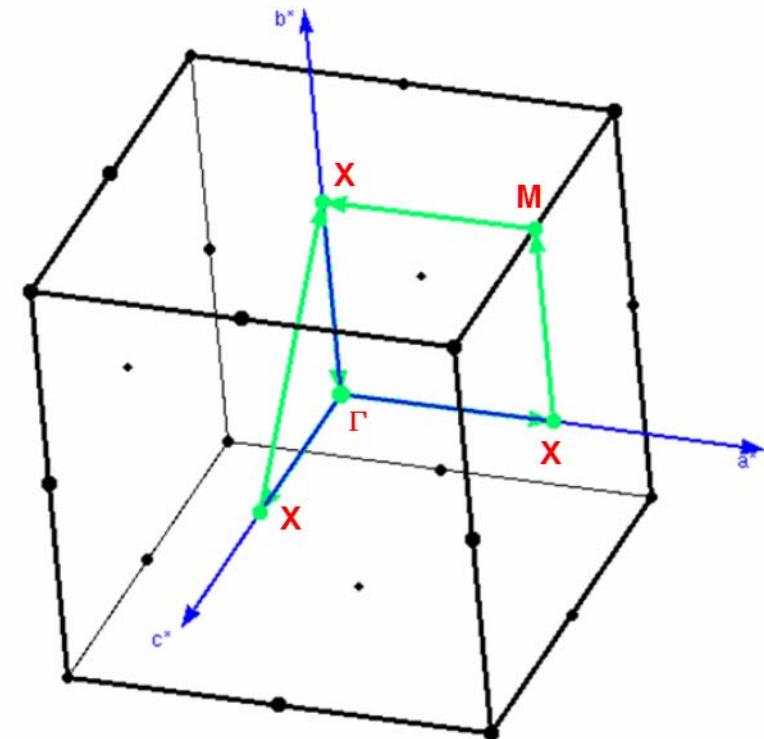
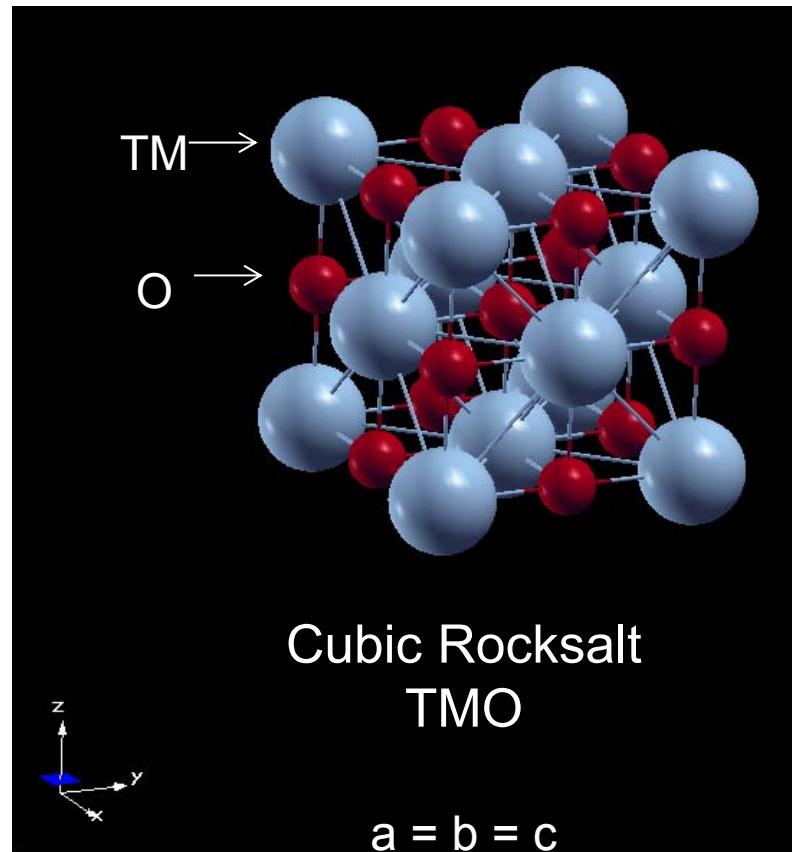
(Charge Transfer Insulators, Instead)



After Imada, et al, RMP 70, 1039 (1998)

Cubic Rocksalt TMOs

Direct and Reciprocal Lattices



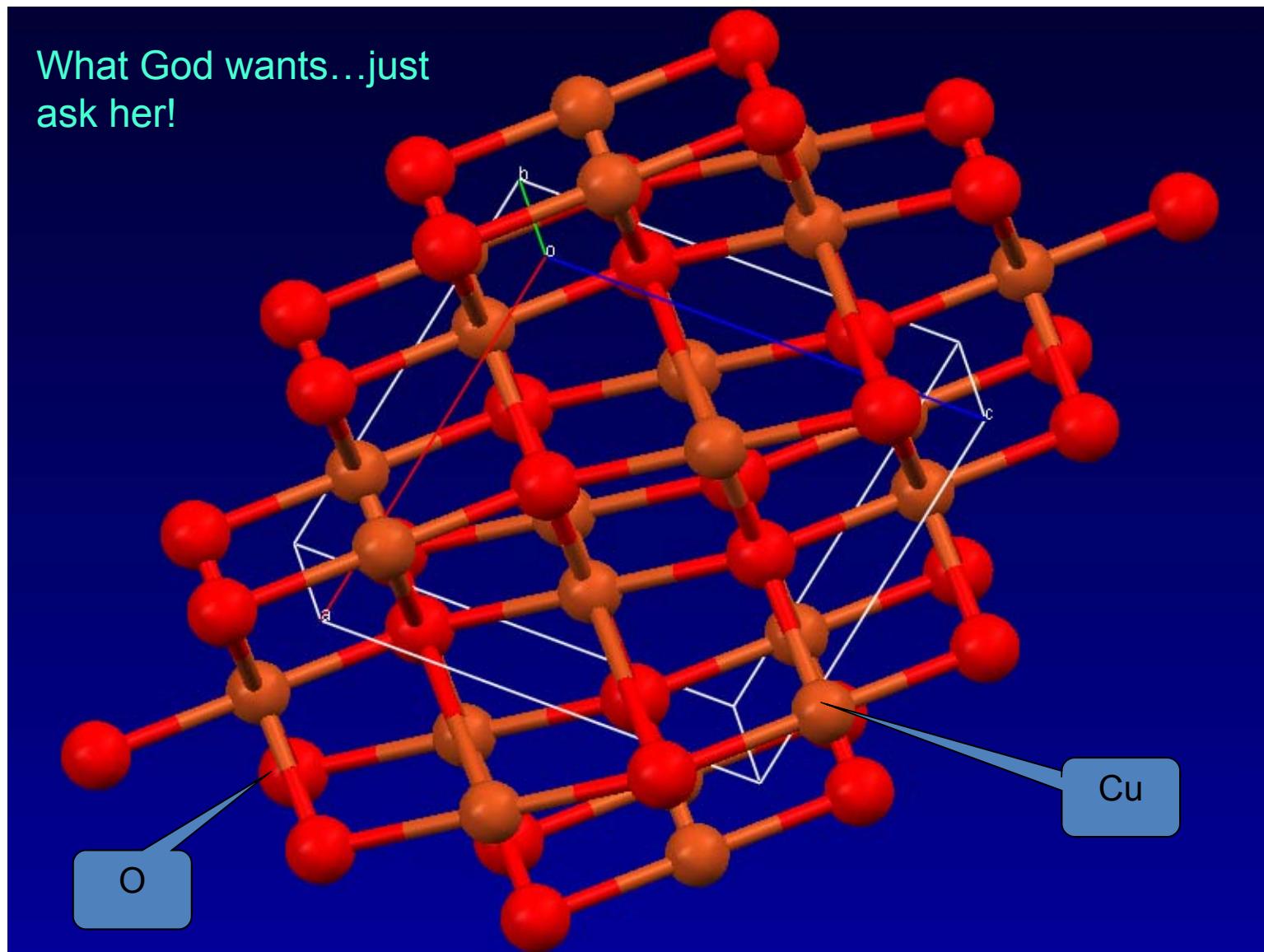
Cubic Rocksalt Divalent TMOs

<u>TMO</u>	<u>3d Config</u>	<u>Properties</u>
MnO	5	MH-CTI (5.6)
FeO	6	MH-CTI (5.9)
CoO	7	MH-CTI (6.3)
NiO	8	MH-CTI (6.5)
CuO	9	XX <i>Doesn't Exist!</i>

See Imada, Fujimori,
Tokura, RPM 70 (1988)

Why Not?

Tenorite (Monoclinic CuO)



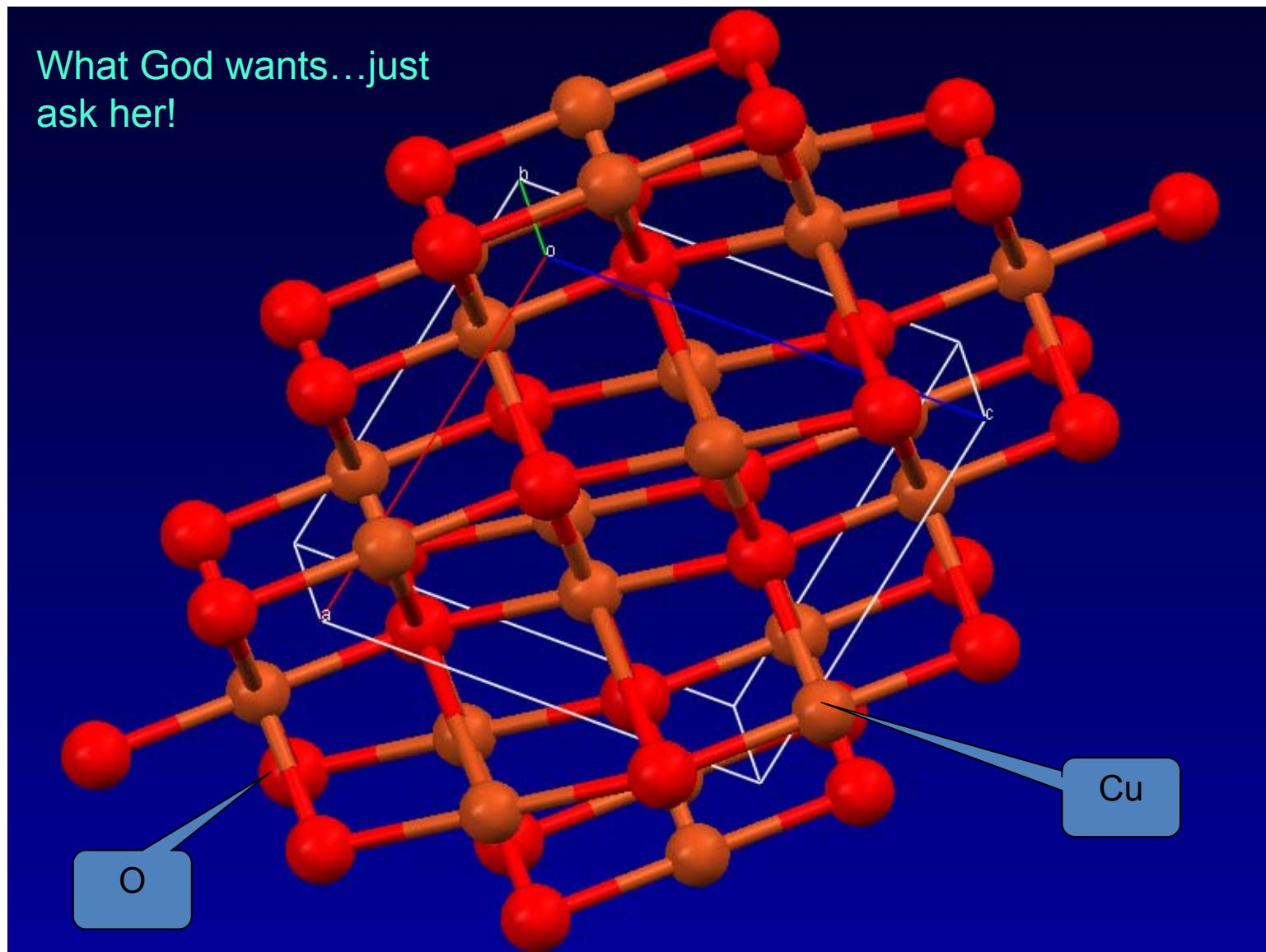
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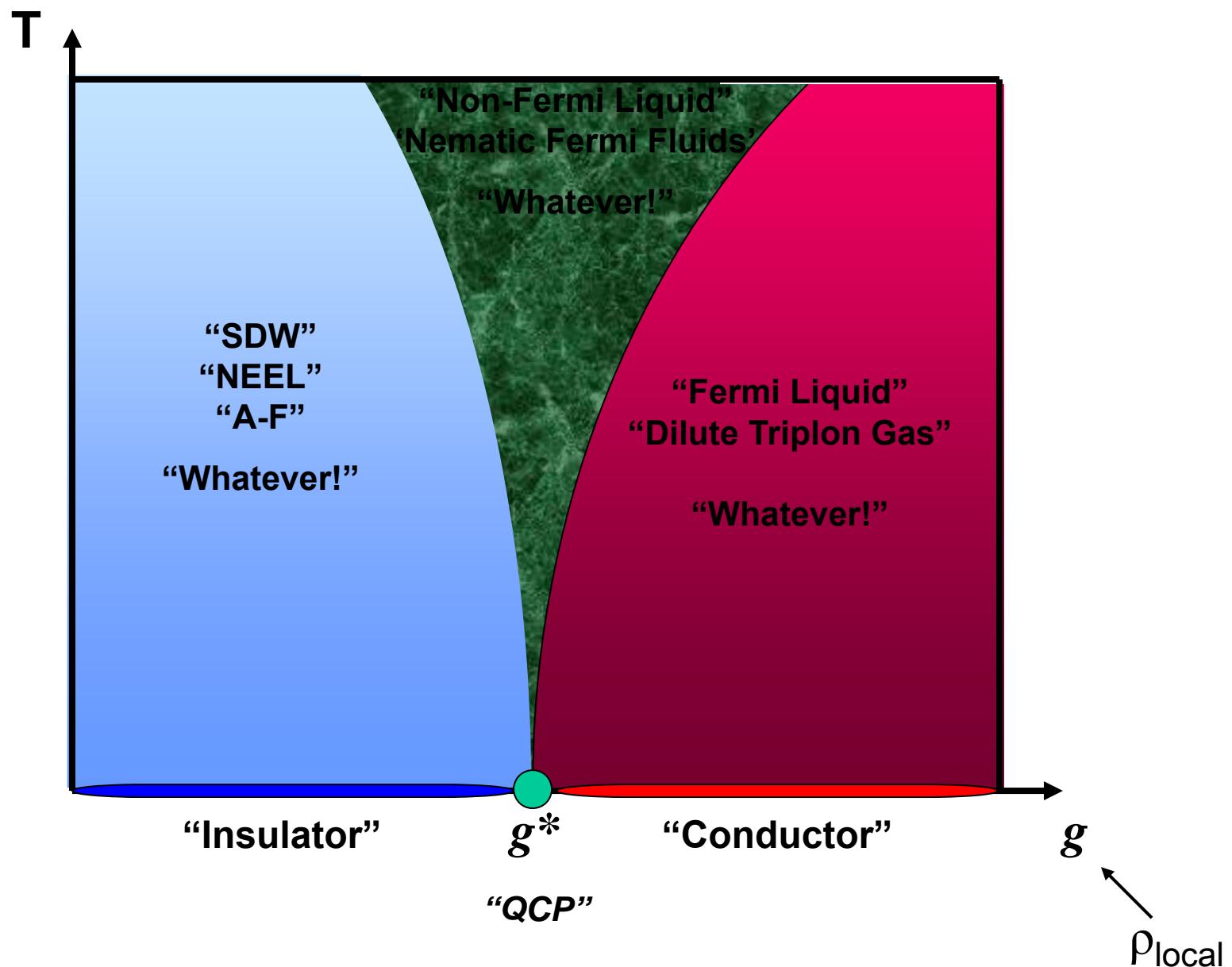
See Imada, Fujimori,
Tokura, RPM 70 (1988)

Why Not?

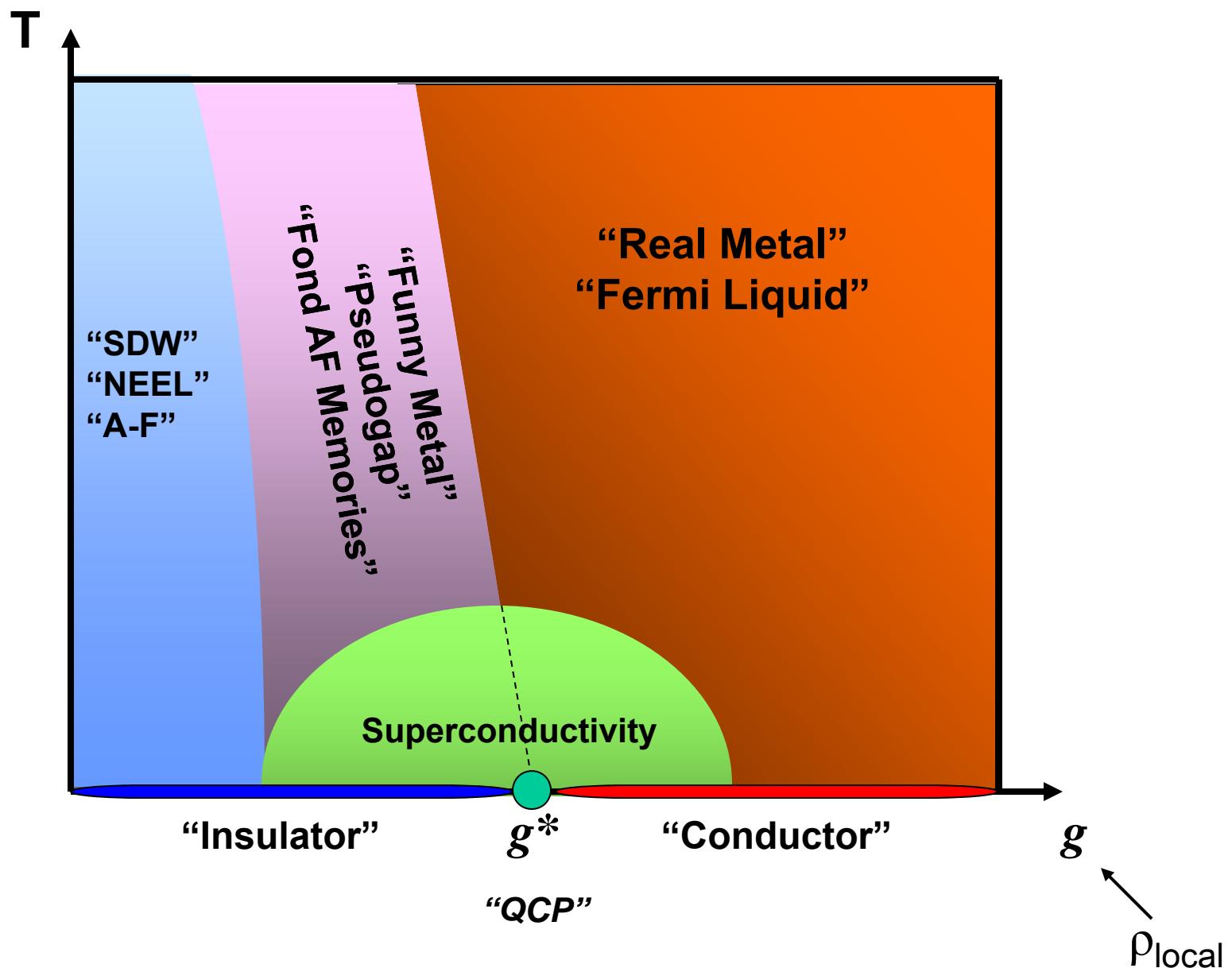
Tenorite (Monoclinic CuO)



The Great Quantum Conundrum



The Colossal Quantum Conundrum



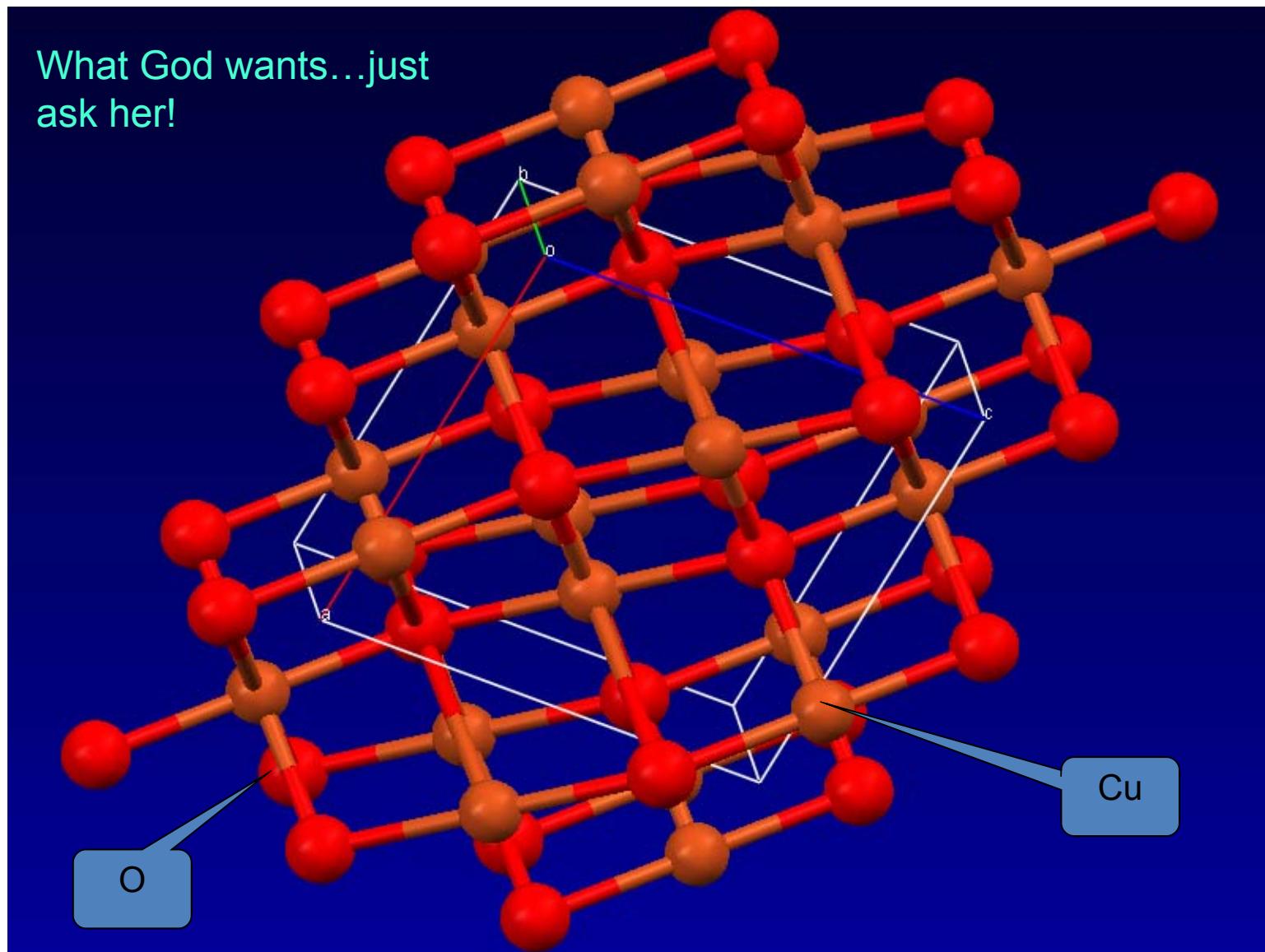
Cubic Rocksalt Divalent TMOs

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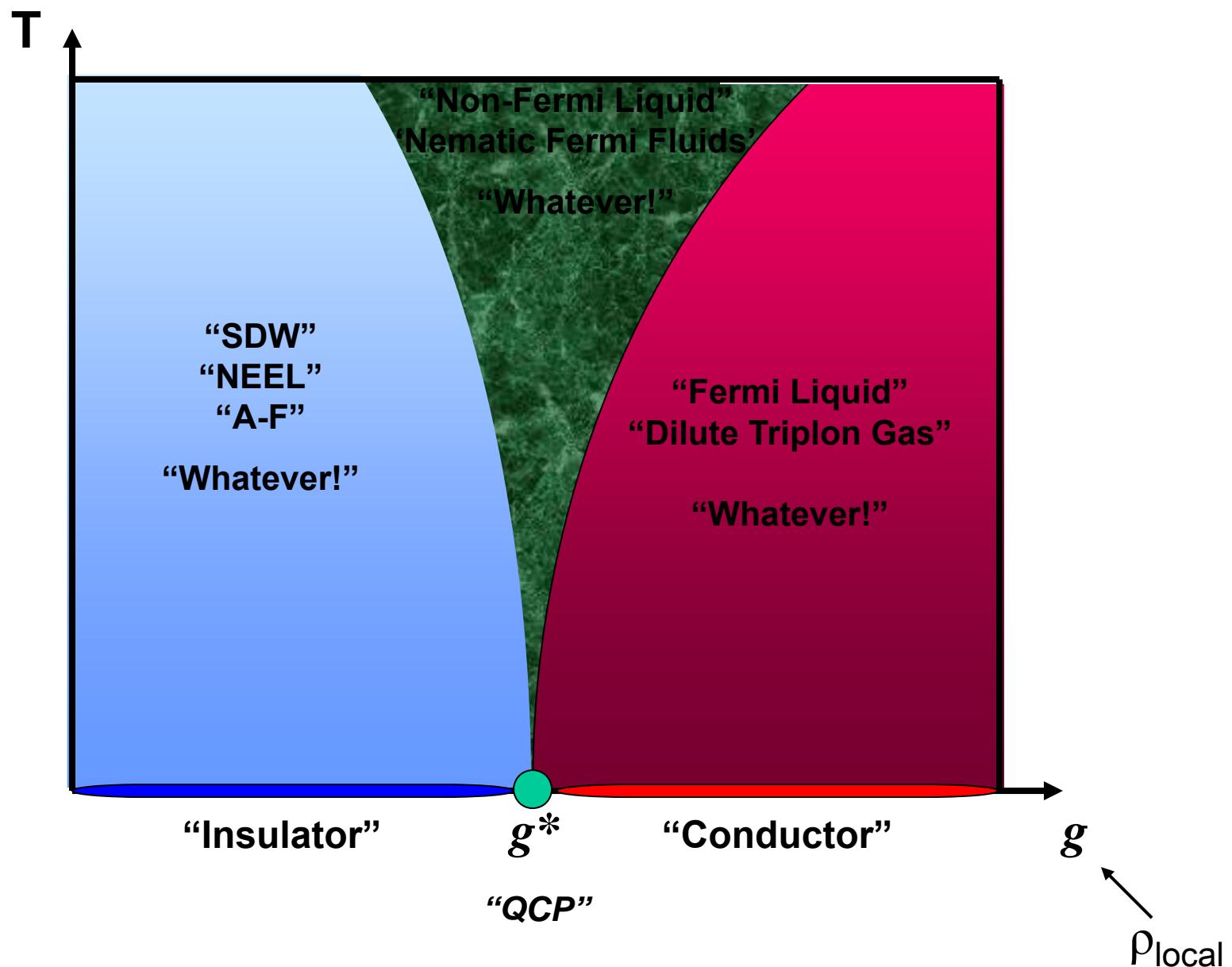
See Imada, Fujimori,
Tokura, RPM 70 (1988)

Why Not?

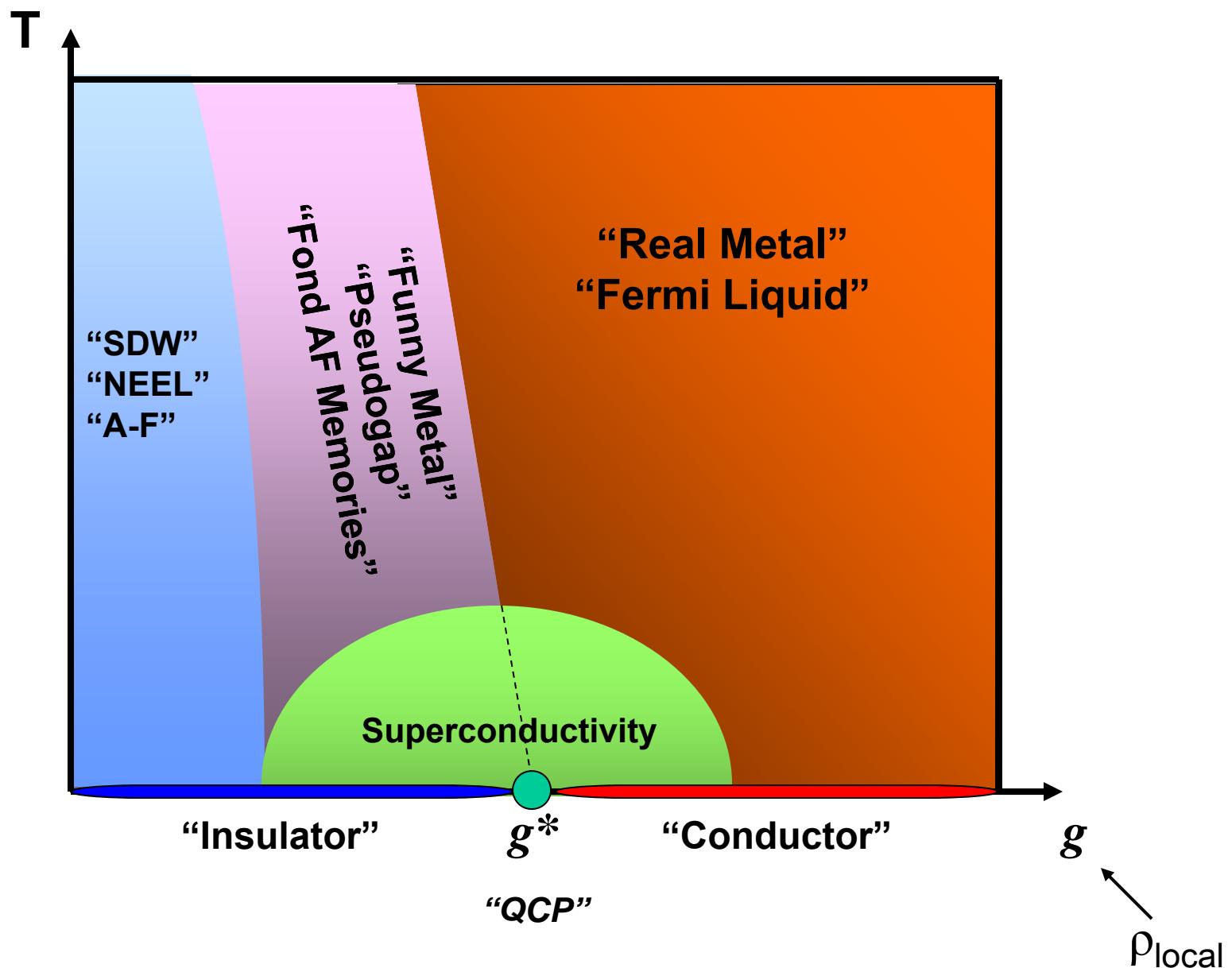
Tenorite (Monoclinic CuO)



The Great Quantum Conundrum



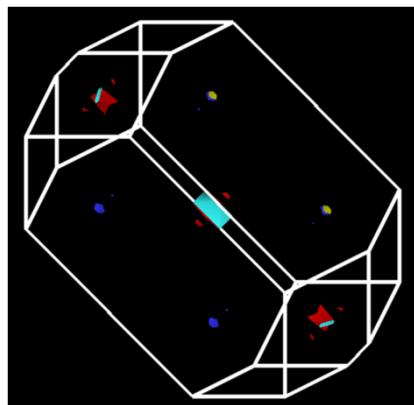
The Colossal Quantum Conundrum



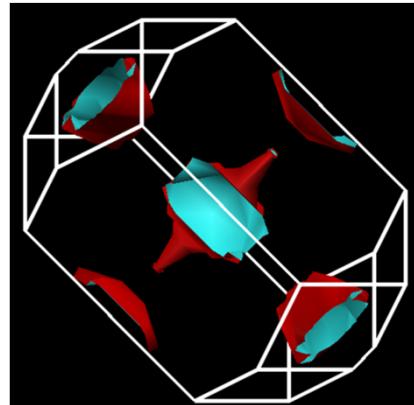
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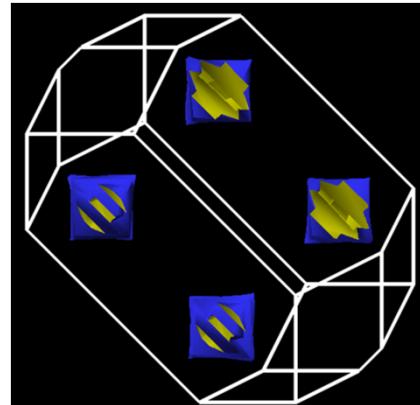
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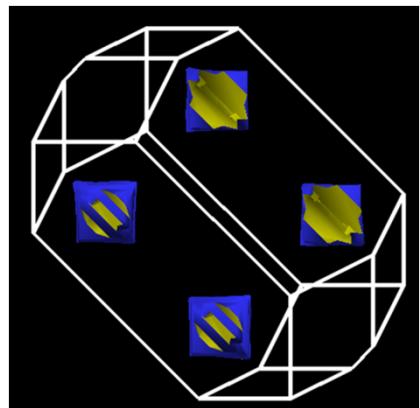
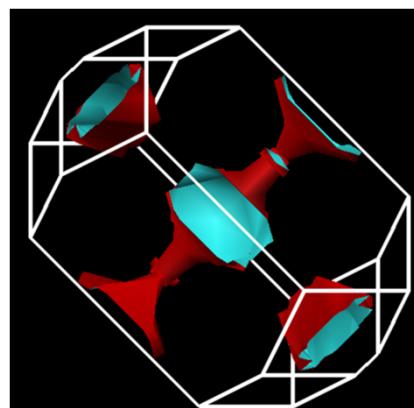
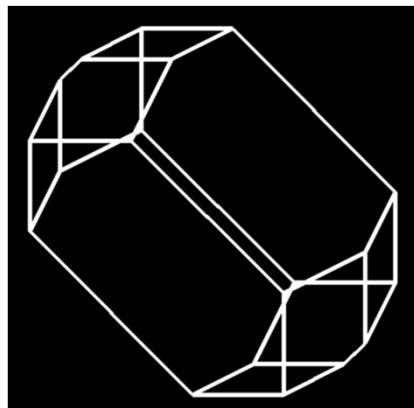
+0.15



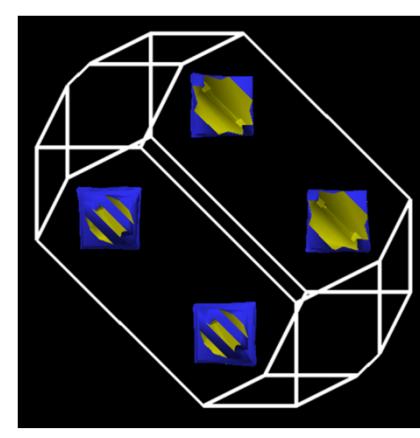
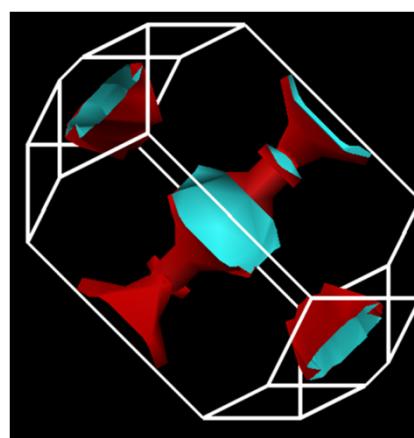
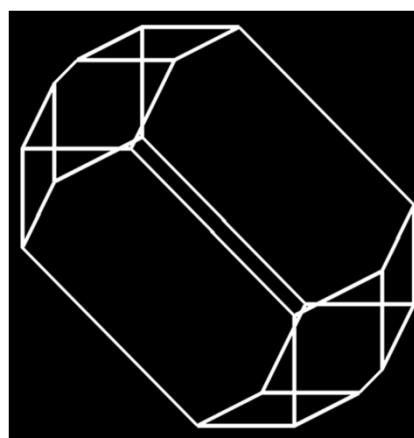
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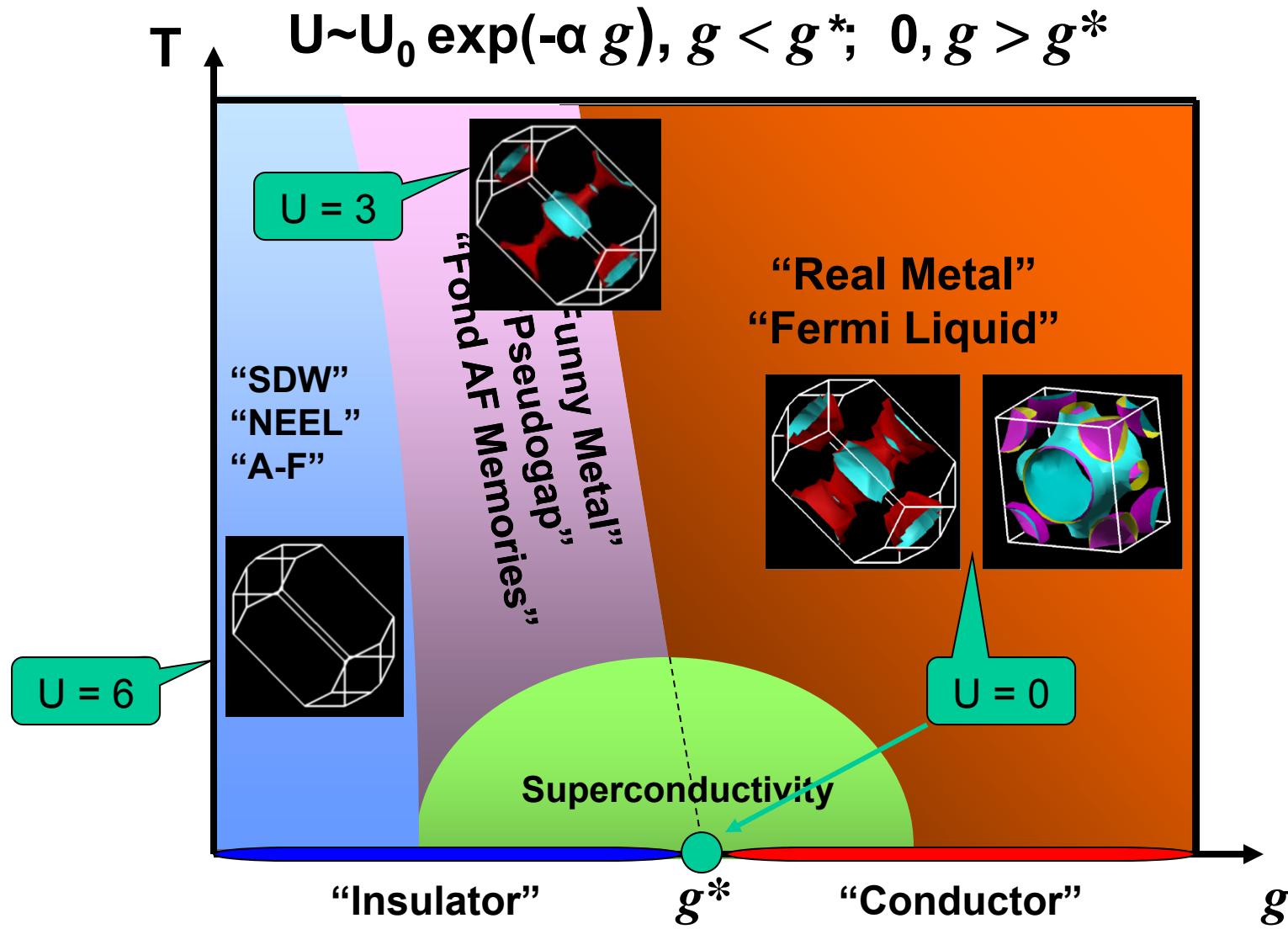
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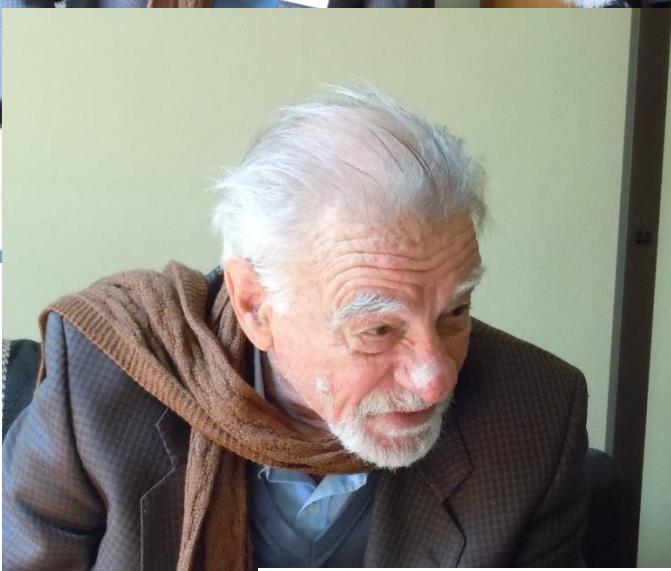
6



The Colossal Quantum Conundrum



Somewhere in here there has to be “BCS-like” pairing!



J Supercond Nov Magn (2011) 24: 1235–1239

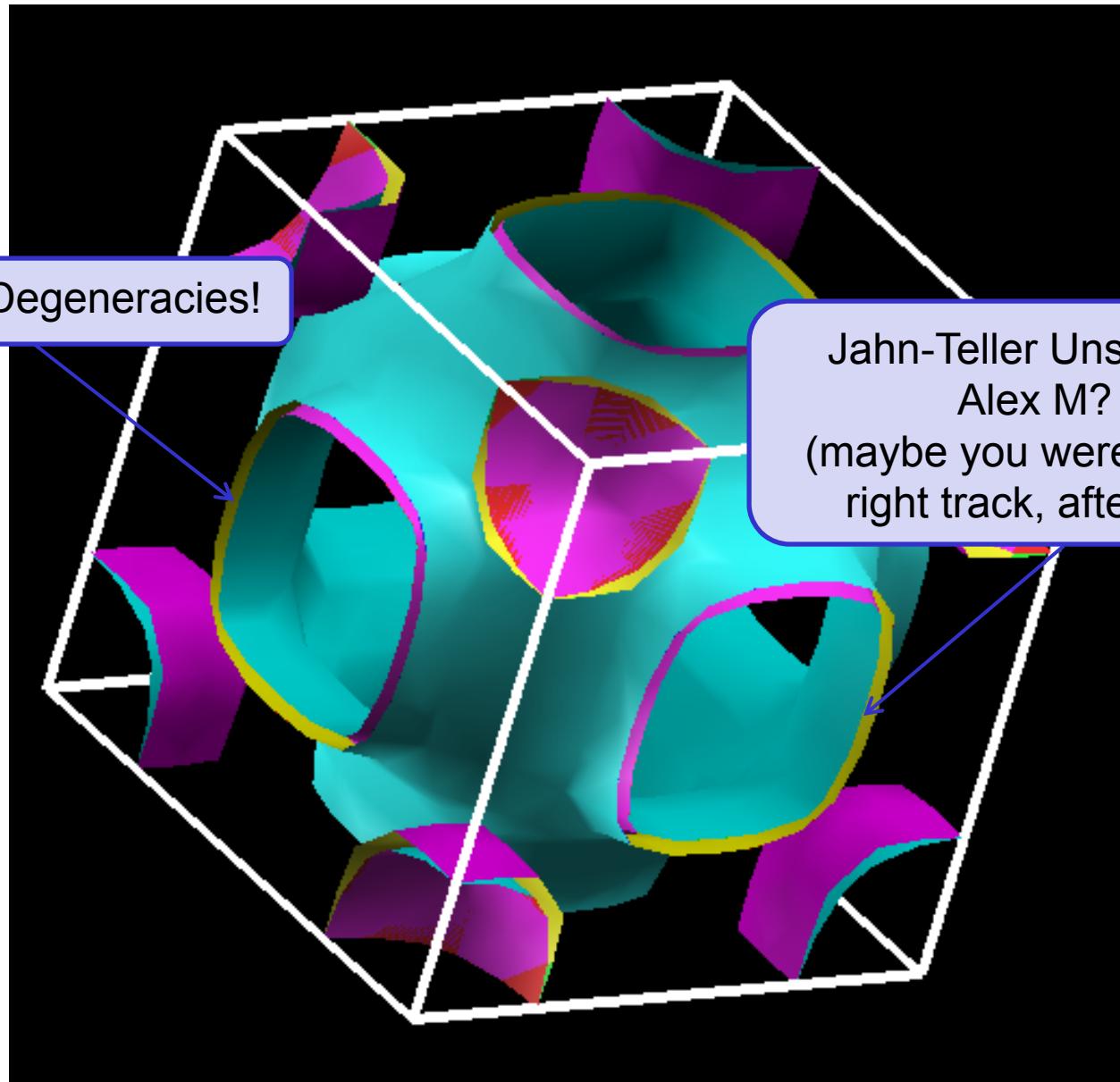
DOI 10.1007/s10948-011-1151-3

**Oxygen Isotope Effect in Cuprates Results from Polaron-induced
Superconductivity**

S. Weyeneth · K.A. Müller

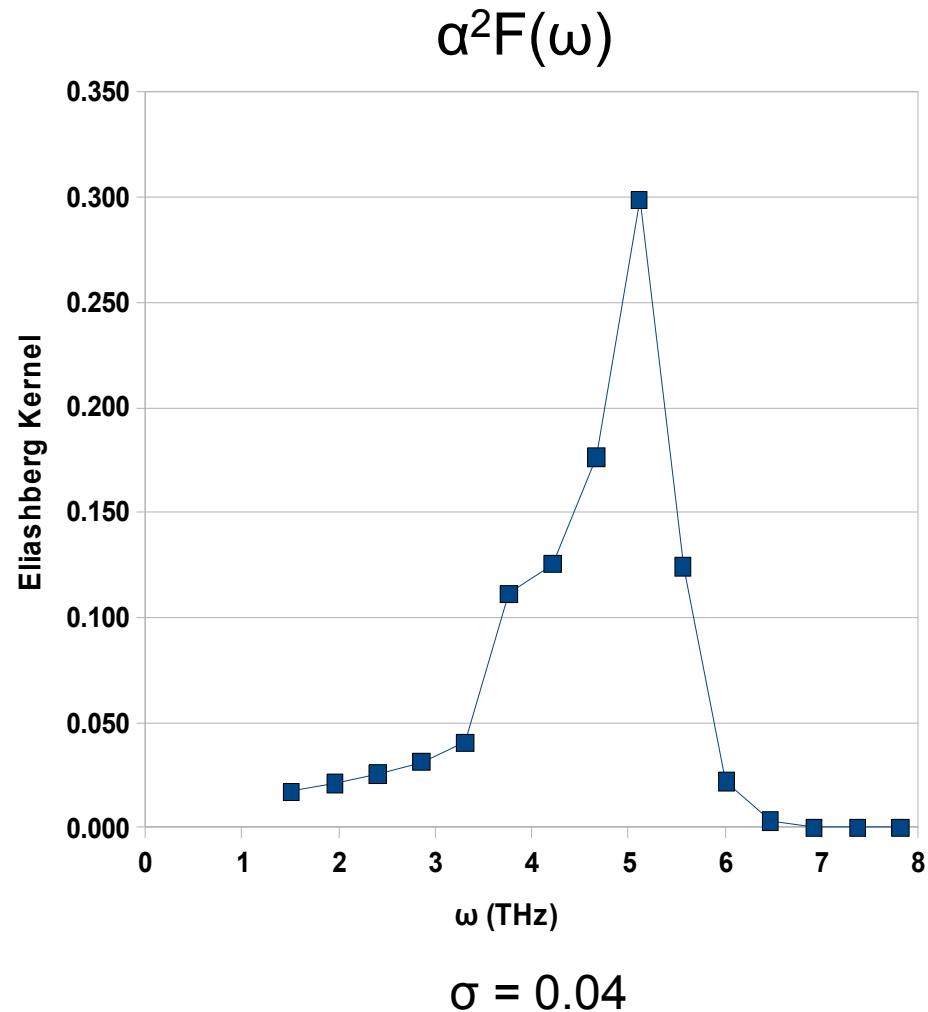


Rocksalt CuO Fermiology ($U = 0$ eV) (8 Bands Combined)



Non-Magnetic ($U = 0$) Cubic Rocksalt CuO

-- Electron-Phonon Properties --



- $\lambda \sim 0.6 - 0.7$
- Consistent with other non-magnetic “HTSCs”

$$T_C = a\Theta e^{-\frac{1}{\lambda - \mu^*}} \quad \lambda k\Theta \ll E_F$$

	T_C (K)	λ	μ^*
K_3C_{60}	16.3	0.51	-
Rb_3C_{60}	30.5	0.61	-
Cs_3C_{60}	47.4	0.72	-

Shakes or Spins or Both?

Are They Copacetic, Competitive...or...

...just another Conundrum?

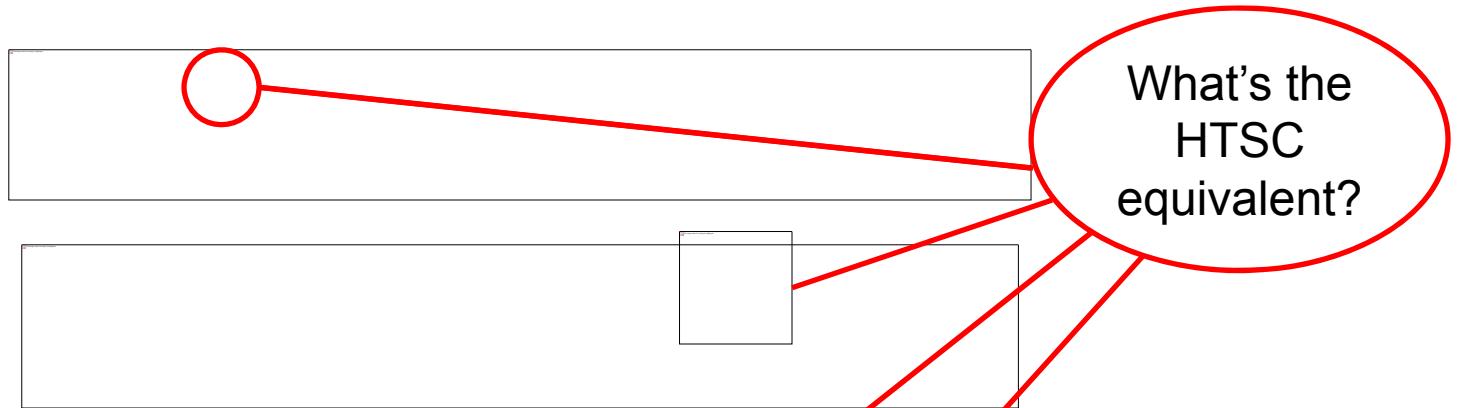
**What formalism is the HTSC analogy to
Migdal-Eliashberg-McMillan?**

(In other words, how do I calculate the value of the BCS gap?)

- Original Strong Coupling, Eliashberg (*JETP*, 1960), *McMillan* (*PR*, 1968)
- Generalized Linhard Response Function (RPA + fluctuations) *Hu and O'Connell* (*PRB* 1989)
- Dielectric Response Function *Kirznits, Maximov, Khomskii* (*JLTP* 1972)

McMillan Strong Coupling

(Computationally implemented by Wierzbowska, et al., cond-mat/0504077, 2006)



$$\alpha^2(\omega) F(\omega) = \int_S \frac{d^2 p}{v_F} \int_{S'} \frac{d^2 p'}{(2\pi\hbar)^3 v_{F'}} \sum_{\nu} g_{pp'\nu}^2 \delta(\omega - \omega_{p-p'\nu}) / \int_S \frac{d^2 p}{v_F}, \quad (19)$$

where the integral $\int d^2 p$ is taken over the Fermi surface and the electron-phonon matrix elements are given by¹⁴

$$g_{pp'\nu} = (\hbar/2MN\sqrt{\omega_{p-p'\nu}})^{1/2} g_\nu(p, p'), \quad (20)$$

where $g_\nu(pp')$ is the electronic matrix element of the change in the crystal potential U as one atom is moved:

$$g_\nu(pp') = \int \psi_p^*(\epsilon_{p-p'}, \nabla U) \psi_{p'} d\mathbf{r}. \quad (21)$$

Generalized Linhard Function

$$\chi^0(\mathbf{q}, \omega) = \sum_{\mathbf{k}, \sigma} \frac{f(\mathbf{k}) - f(\mathbf{k} + \mathbf{q})}{\hbar\omega - (\varepsilon_{\mathbf{k} + \mathbf{q}} - \varepsilon_{\mathbf{k}}) + iDq^2}$$

$$D = \lim_{t \rightarrow \infty} \frac{1}{2t} \overline{\delta \mathbf{R}^2(t)}$$

$$V(\mathbf{q}) = 4\pi e^2/q^2$$

$$\epsilon(\mathbf{q}, \omega) = 1 - V(\mathbf{q})\chi^0(\mathbf{q}, \omega)$$

$$\begin{aligned} \epsilon_1(x, y) = & 1 + \frac{q_{TF}^2}{8k_F^2 x^2} \left[1 + \frac{1}{8x} \left[(1 + b^2 x^2 - \nu_+^2) \ln \left[\frac{(1 + \nu_+)^2 + b^2 x^2}{(1 - \nu_+)^2 + b^2 x^2} \right] + (1 + b^2 x^2 - \nu_-^2) \ln \left[\frac{(1 + \nu_-)^2 + b^2 x^2}{(1 - \nu_-)^2 + b^2 x^2} \right] \right] \right. \\ & \left. - \frac{b}{2} \left\{ \nu_+ \left[\arctan \left[\frac{1 - \nu_+}{bx} \right] + \arctan \left[\frac{1 + \nu_+}{bx} \right] \right] + \nu_- \left[\arctan \left[\frac{1 - \nu_-}{bx} \right] + \arctan \left[\frac{1 + \nu_-}{bx} \right] \right\} \right] \end{aligned}$$

where

$$x = \frac{q}{2k_F}, \quad y = \frac{\hbar\omega}{4\varepsilon_F}, \quad q_{TF}^2 = \frac{4me^2k_F}{\pi\hbar^2}, \quad b = \frac{2mD}{\hbar}, \quad \nu_{\pm} = x \pm y/x,$$

“Fluctuations?”
“Empirical?”

HO (1989)

Dielectric Response Function

$$G(\mathbf{k}, i\omega_n) = 1/(i\omega_n - \xi_k)$$

$$\begin{aligned} F(\mathbf{p}, i\omega_n) = & -G(\mathbf{p}, i\omega_n)G(-\mathbf{p}, -i\omega_n)T_c \sum_m \int [d^3k/(2\pi)^3] \\ & \times V(\mathbf{p} - \mathbf{k}, i\omega_n - i\omega_m)F(\mathbf{k}, i\omega_m) \end{aligned}$$

$$V(\mathbf{q}, i\omega_n) = \frac{4\pi e^2}{q^2} \left[1 - \int_0^\infty \frac{dE^2 \rho(\mathbf{q}, E)}{\omega_n^2 + E^2} \right]$$

In principle, KMK can calculate the BCS gap for general “bosonic” fields, be they phonons, magnons, spin-ons, excitons, plasmons...or morons!

So...

- Let's "Shut up and start calculating."
 - David Mermin, Cornell, as quoted by yours truly, in [Nature 4 August 2011](#).
- I'd like to help...
 - Then,

Bottom Line

Can studying CuO proxies with DFT
+ LDA+U
+ phonons
+ spins
provide insight into the origins of High-T_C?

I say “Yes,” but...
As Bob Laughlin says, “Size Matters...!”
...and I need a...
BIGGER COMPUTER!