

A DFT Study of Tetragonal Rocksalt Proxy Copper Monochalcogenide Structures: -- Implications for Possible High-Tc Superconductivity --



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Paper 8 **9:24 AM – 9:36 AM** Mile High Ballroom 4F

Session Y47 **Theory of Strongly Correlated Superconductivity** 8 AM – 11 AM, Friday, 7 March

- Our Computational Tool Box -

- DFT + Hubbard U
 - Quantum Espresso
 - Bands, Fermiologies, States (DOS), Phonons
- Graphics
 - Xcrysden, XMGRACE
 - Bandwidths, Fermi Surfaces, Projected DOS
- Modeling
 - Neel Temperatures a la Van Vleck/Anderson/Hubbard
 - Superconductivity via Eliashberg/McMillan

The Various Flavors of Copper "Monoxide"



"Configuration/Coordination Space"

Rocksalt af-CuX Crystallography

nm-Translational Unit Cell

af-Primitive Cell



two Cu ions to correctly represent the doubly-periodic spin up/down ordering



CuX (Cubic, Equilibrium Lattice Constant(s))

CuSe (a ≈ 5.0 Å)



CuTe (a ≈ 5.3 Å)



What Does Experiment Say About Rocksalt CuO?

It's <u>*Tetragonal(!)*</u> for 4-6 monolayers forced-epi grown on STO yielding a film with lattice constants a = b = 3.905 Å, and $c/a \approx 1.3$, representing a <u>5% basal-plane</u> <u>contraction</u> down from pure cubic having a = b = c = 4.1 Å. (Siemons, et al, PRB 79, 195122 (2009))

CuX (Tetragonal)

(Assuming a 5% contraction of the a, b lattice constants a la CuO on STO)

CuSe (a = b \approx 4.75 Å; c/a \approx 1.1)

CuTe (a = b ≈ 5.05 Å; c/a ≈ 1.1)

CuSe

CuSe

CuSe

CuSe

Van Vleck /Anderson/Hubbard Model of Neél Temperature

Take the definition of "Exchange Energy" from Anderson (1959):

$$J_{ij} = 2t_{ij}^2 / U$$

where t_{ij} is the transfer integral from the spin states on one TM ion, directly to a neighbor, or through an intervening anion, and U is the on-site Hubbard coulomb repulsion potential, e.g.,

Now, "plug into" Van Vleck (1938-42), within the "molecular field approximation," to get T_N :

$$T_{N} = \frac{2S(S+1)}{3k_{B}} \sum_{i \neq j} J_{ij} = \frac{4S(S+1)}{3k_{B}U} \sum_{i \neq j} t_{i \neq j}^{2}$$

Here S is the net cationic spin and k_B the Boltzmann constant (8.61733•10⁻⁵ eV/°K). The "transfer integrals" are given by,

$$t_{ij} = \left\langle \phi_j \left| H \left| \phi_i \right\rangle \approx "Bandwidth" / 4 = w_{i,j} / 4 \right. \right.$$

where the ¢'s are the "spin carrying" orbitals and H a "tight-binding-like Hamiltonian."

Néel Temperature vs. TMO

Néel Temperature vs. TMO

Superconductivity(?) Phonons(?) Eliashberg-McMillan

$$H_{el-ph} = \sum_{\mathbf{k},\mathbf{q},\nu} g_{\mathbf{k}+\mathbf{q},\mathbf{k}}^{\mathbf{q}_{\nu,mn}} c_{\mathbf{k}+\mathbf{q}}^{\dagger m} c_{\mathbf{k}}^{n} (b_{-\mathbf{q},\nu}^{\dagger} + b_{\mathbf{q},\nu})$$

$$\lambda_{\mathbf{q},\nu} = \frac{2}{N(\varepsilon_F)\omega_{\mathbf{q},\nu}} \sum_{mn} \sum_{\mathbf{k}} \left| g_{\mathbf{k}+\mathbf{q},\mathbf{k}}^{\mathbf{q}_{\nu,mn}} \right|^2 \delta(\varepsilon_{\mathbf{k}+\mathbf{q},m} - \varepsilon_F) \delta(\varepsilon_{\mathbf{k},n} - \varepsilon_F)$$

$$\alpha^{2}F(\omega) = \frac{1}{N(\varepsilon_{F})} \sum_{mn} \sum_{\mathbf{q},\nu} \delta(\omega - \omega_{\mathbf{q},\nu}) \sum_{\mathbf{k}} \left| g_{\mathbf{k}+\mathbf{q},\mathbf{k}}^{\mathbf{q}_{\nu,mn}} \right|^{2} \delta(\varepsilon_{\mathbf{k}+\mathbf{q},m} - \varepsilon_{F}) \delta(\varepsilon_{\mathbf{k},n} - \varepsilon_{F})$$

$$\lambda = 2 \int_0^\infty \frac{\alpha^2 F(\omega)}{\omega} d\omega = \sum_{\mathbf{q},\nu} \lambda_{\mathbf{q},\nu}$$

Need to compute $g_{\mathbf{k}+\mathbf{q},\mathbf{k}}^{\mathbf{q}_{\nu,mn}}$!

e-p Interaction in the DFT/LDA Formalism

$$g_{\mathbf{k}+\mathbf{q},\mathbf{k}}^{\mathbf{q}_{\nu,mn}} = \sqrt{\hbar/2\omega_{\mathbf{q},\nu}} \left\langle \psi_{\mathbf{k}+\mathbf{q},m} \left| \Delta V_{KS}^{\mathbf{q},\nu} \right| \psi_{\mathbf{k},n} \right\rangle$$

$$\Delta V_{KS}^{\mathbf{q},\nu} = \sum_{\mathbf{R}} \sum_{s} \frac{\partial V_{KS}}{\partial \vec{u}_{s,\mathbf{R}}} \cdot \vec{u}_{s}^{\mathbf{q},\nu} \frac{e^{i\mathbf{q}\cdot\mathbf{R}}}{\sqrt{N}}$$

$$T_{C} = \frac{\Theta_{D}}{1.45} \exp\left(-\frac{1.04(1+\lambda)}{\lambda - \mu^{*}(1+0.62\lambda)}\right)$$

The Colossal Quantum Conundrum

Perhaps phonon-mediated?

Well, how about the "U = 0, Fermi Liquid" limit for doped proxy tet-CuO?

So let's do it and "compute" what happens!

q = 0.15 |e|/CuO (holes)

≈ 43 °K

q = -0.15 |e|/CuO (electrons)

≈ 25 °K

Apply DFT to obtain $g_{k+q,k}^{q_{\nu,mn}}$ between electrons and phonons, followed by application of the Eliashberg-McMillan-Allen-Dynes formalism to find Tc:

Can We Really Make Any of This Stuff?

Forced-epitaxial thin film growth is obvious choice (as it was with tet-CuO. Substrate selection likely limited, but here are possible choices:

- 1. CuS (4.7 Å)Rocksalt ZnO (4.580 Å, ~3% compression)
Rutile TiO2 (4.591 Å, ~2.5% compression)2. CuSe (5.0 Å)Hex $Al_2O_3 (4.748 Å, ~5\% compression)$ 3. CuTe (5.3 Å)Cubic $ZrO_2 (5.147 Å, ~5.3\% compression)$
 - YSZ (5.13 5.23 Å, ~3.5% compression max)

CaF (5.46 Å, ~3% expansion)

Methodologies

- a) MBE PLD:
 - i. Use appropriate sintered sample source.
 - ii. Empirically determine optimum substrate temperature and argon pressure.
 - iii. Characterize growth and structure via in-situ "high pressure compatible" RHEED, XPS, UPS, LEED.
- b) External characterization, depending on stability in air:
 - i. 4-probe transport.
 - ii. UV-Vis optical transmission and reflectivity.

"An Ideal Lab"

The Bottom Line(s)

- For X = S, Se and Te, neither a finite U or a "5% basal" tetragonal distortion has much effect on their respective CuX Fermiologies, and likely transport/magnetic properties dependent thereon.
- However, the respective Fermi surfaces ...may...may... contain nesting topologies promoting itinerant antiferromagnetism a la Cr, but, unlike Cr, here for X = S, Se, Te, the DOS at Ef is dominated by p-like chalcogenide overlap.
- Future homework for proxy structure modelling, suggested by preliminary results on "doped" tet-CuO: Let's look for electron-phonon mediated superconductivity!
- But ...most importantly... experiment *always* rules. Our fundamental computational finding is that equilibrium rocksalt CuS, CuSe and CuTe structures can in principle exist ...so let's try to make and dope them and henceforth measure their properties!

Finally, there is something quite special about the Cu-O bond in squareplanar symmetry!

...but we knew that already... in 1986 B & M told us so!