

Progress in High Temperature Superconductivity – Vol. 31

Proceedings of the XII Winter Meeting on Low Temperature Physics

# **SUPERCONDUCTING CERAMICS**

Vista Hermosa, Morelos, México 13 – 16 January 1991

**Editors**

**J. L. Heiras**

**L. E. Sansores**

**A. A. Valladares**

**World Scientific**

Progress in High Temperature Superconductivity – Vol. 31

---

**XII Winter Meeting on Low Temperature Physics**

# **SUPERCONDUCTING CERAMICS**

Vista Hermosa, Morelos, México 13 – 16 January 1991

**Editors**

**J. L. Heiras**

**L. E. Sansores**

**A. A. Valladares**

*Instituto de Investigaciones en Materiales  
Universidad Nacional Autónoma de México*



**World Scientific**

*Singapore • New Jersey • London • Hong Kong*

## CONTENTS

Preface	vii
BiSrCaCuO High $T_c$ Superconductors	1
<i>O. Laborde, P. Monceau, H. Raffy &amp; S. Labdi</i>	
Metalorganic Deposition of High- $J_c$ $Ba_2YCu_3O_{7-x}$ Thin Films on Single Crystal Substrates	13
<i>M. J. Cima, P. C. McIntyre, J. A. Smith, Jr., B. P. Chang &amp; R. B. Hallock</i>	
Critical Currents in Proton Irradiated $YBa_2Cu_3O_{7-\delta}$ Crystals	25
<i>L. Civale, M. W. McElfresh, A. D. Marwick, T. K. Worthington, A. P. Malozemoff, F. Holtzberg, C. Field, J. R. Thompson, D. K. Christen &amp; M. A. Kirk</i>	
Electron Tunneling in Superconducting Ba-K-Bi-O. A Review	45
<i>R. Escudero</i>	
Thermodynamic Considerations in the Y-Ba-Cu-O Systems	55
<i>R. Beyers &amp; B. T. Ahn</i>	
Review of Organic Conductors: A Structural Analysis	85
<i>R. R. Schumaker</i>	
Stepped Copper Surfaces Observed by Scanning Tunneling Microscopy	121
<i>S. Rousset, S. Gauthier, O. Siboulet, W. Sacks, M. Belin, J. Klein &amp; M. A. Huerta Garnica</i>	
Experimental Predictions of the Hole Mechanism of Superconductivity	129
<i>F. Marsiglio</i>	
First Principles Cluster Investigation of the Electronic Structures of Copper Oxide Superconductors	141
<i>M. Eto &amp; H. Kamimura</i>	
Oxygen Ordering in $YBa_2Cu_3O_{6+\delta}$	165
<i>J. L. Morán-López &amp; J. M. Sanchez</i>	

The Importance of Being <i>N</i> -type <i>P. M. Grant</i>	175
Author Index	185

## THE IMPORTANCE OF BEING N-TYPE

P. M. GRANT\*

*Instituto de Investigaciones en Materiales, Universidad Nacional Autónoma de México, Apartado Postal 70-360, 04510 México, D.F.*

## ABSTRACT

Perhaps the most significant recent advance in high-Tc materials has been the discovery of superconductivity in the  $\text{Nd}_2\text{CuO}_4$ -based systems. Certain solid state measurements, primarily Hall effect and thermoelectric power, would suggest transport in the normal state, and by implication also the superconducting state, of these new compounds is dominated by electrons rather than holes, the latter having been the case for all previous high-Tc materials. This discovery of charge reciprocity is potentially important for guiding theoretical developments. However, the n-type compounds differ significantly in both their preparation and structural properties from their p-type antecedents, and display important non-reciprocal behavior in doping and oxidation-reduction treatments. In this article, we will review the present situation regarding n-type high temperature superconductors, paying particular attention to the aforementioned differences from the hole carrier materials and the impact such differences may have on extending the number of presently known n-type systems.

\* On sabbatical from the IBM Almaden Research Center, San Jose, California, USA.

Electron-hole duality has long been a mainstay in the modern theory of the electronic properties of semiconductors and metals. This was not always the case. In the 19th century, it was very perplexing to discover that some metals had a positive Hall coefficient, and even worse, many yielded Seebeck coefficients opposite in sign to that of the Hall effect. These questions were resolved in large part with the arrival of the quantum theory of solids in the 1930s and 40s, although some, especially those involving the Seebeck effect in noble metals, still seek complete answers today.

As regards superconductivity, conventional BCS theory makes no distinction between holes and electrons. Indeed, the number of superconductors with a hole-like normal state is manifold, and there is some feeling (see the paper by Marsiglio in this conference [1], and the one presented by Hirsch at a preceding meeting in this series held in Cocoyoc [2]) that this may be a universal property of all superconductors, requiring, perhaps, extensive departures from the BCS concept.

All early high temperature superconductors displayed the usual aspects of possessing positive charge normal state carriers...positive Hall and Seebeck response, reinforced by simply adding up the excess charge expected from conventional ionic charge states on the constituent ions. From the very beginning attempts were made (all unsuccessful, and none, to the knowledge of the author, ever published) to discover electron analogs to the  $\text{La}_2\text{CuO}_4$  system. Those participated in by the author included tetravalent ion substitutions (e.g., Pb and Sn) and the use of various vanadate hosts (the electron equivalent of  $\text{Cu}^{++}$ ). In fact, the literature already included several strong hints as to how to accomplish this objective. A paper by Shaplygin, et al. [3], preceding by

integrations of dyad products of the group velocity (i.e., dispersion in the energy vs. quasimomentum band structure) over the Fermi surface which in principle could be dominated by small regions of high curvature. This situation is unlikely given the small carrier concentration in high- $T_c$  superconductors relative to most other metals, so that the carriers are very likely confined to band extrema where the effective mass approximation holds and an unambiguous sign can be attached to the carriers. More uncertain is whether band overlap occurs, resulting in electron-hole pockets in which electrons dominate the normal state, yet holes still form the paired state. If, in the p-type compounds, holes reside in O 2p - Cu 3d hybridized valence bands, and, for the n-type case the electrons transport in the Cu 4s conduction band, the most consistent interpretation of core level spectra obtained for each material class [9], then band overlap is improbable given the usual magnitude of the energy gap in transition metal oxides between these two bands. Thus it is prudent to conclude at the present time that charge reciprocity in high temperature superconductors exists and must be accommodated by any theory put forward to describe these materials.

The existence of charge symmetry does not imply structural or chemical symmetry between n- and p-type high- $T_c$  compounds such as found, for example, in covalent semiconductors. Important differences in both crystallographic structure and chemical preparation arise. In Fig. 1 we compare unit cells of the two compounds and the respective doping sites. The principal difference is the coordination of the lanthanide cation with respect to the nearest neighbor oxygen anions. In  $\text{La}_2\text{CuO}_4$ , there is Cu-O plane apical oxygen which in  $\text{Nd}_2\text{CuO}_4$  moves to its observed position on the unit cell face from its former edge location. Apparently this lack of an apical oxygen is the key structural feature which permits electron delocalization. The host structure choice

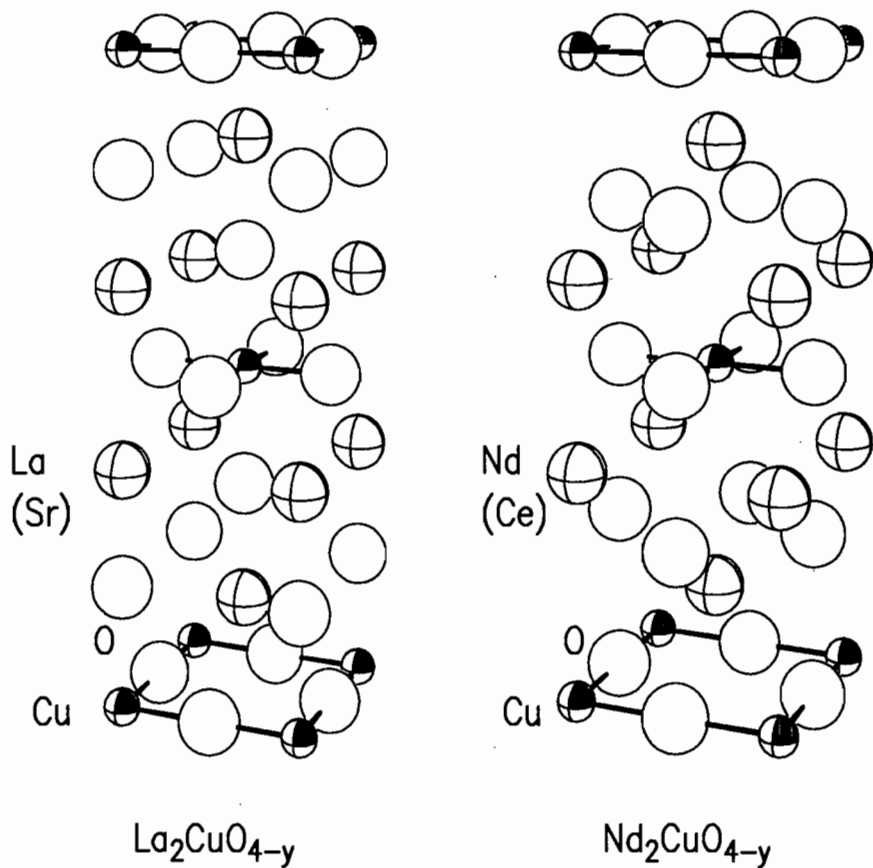


Fig. 1. Unit cells of  $\text{La}_2\text{CuO}_{4-y}$  and  $\text{Nd}_2\text{CuO}_{4-y}$ .  
Note the absence of apical oxygens  
in the latter structure.



possessing this feature for potential n-type superconductivity has been largely restricted to the  $\text{Nd}_2\text{CuO}_4$  class, inasmuch as all other Cu-O planar perovskite structures have at least one apically coordinated plane and appear p-type, with any accompanying non-apical planes "dead" with regard to transport; e. g., the middle Cu-O layer in the 2223 structure. As has been pointed out [10], the single exception was the "infinite layer" compound  $\text{Ca}_{0.86}\text{Sr}_{0.14}\text{CuO}_2$ , which was discovered by Siegrist [11] and coworkers in mid-1988. Recently, two structural isomorphs,  $\text{Sr}_{0.86}(\text{Nd,Pr})_{0.14}\text{CuO}_2$ , were synthesized by Smith, et al., [12] and both were found to have superconducting transitions near 40K. By unit cell charge count, the substitution of trivalent Nd or Pr should result in excess negative charge and thus an n-type normal state. However, it was found by the authors of ref. [12] that their preliminary measurements of the Hall constant were positive, a finding they speculate is due to minority phases in their samples. If subsequent research reveals  $R_H < 0$ , especially near the transition temperature, as additional confirmation of charge counting, then the discovery of n-type superconductivity at 40K is an important advance in at least two aspects. Firstly, unlike the  $\text{Nd}_2\text{CuO}_4$ -based systems, the "infinite layer" compounds do not require direct reduction processing. On the contrary, both the precursor compounds and the final phase are obtained via high pressure oxygen processing. However, when more is known about the microstructure, it may turn out that oxygen vacancies indeed exist. Secondly, and most critically, the discovery of this new high-Tc material demonstrates that transition temperatures as high as those obtained in Sr-doped  $\text{La}_2\text{CuO}_4$  can be reached in systems without apical oxygens; thus proving that the existence of this ion is not a necessary condition for truly high temperature superconductivity.

Where do we go from here? Superconductivity at 60K has been reported in a "double layer" variant of doped- $\text{La}_2\text{CuO}_4$  [13]. In

fact, a whole series of multiple layer variants seem possible [14]. A suggestive path would be to attempt the synthesis of appropriately n-doped T' isomorphs where trivalent rare earth replacement of the lanthanum would move the apical oxygen to the cell edges as in  $\text{Nd}_2\text{CuO}_4$ . Only time and clever chemistry will tell.

The astute reader may recognize the title of this article as a pun, albeit far-fetched, on Oscar Wilde's famous 1895 play, "The Importance of Being Earnest." Elgernon, who originally pretended to be Jack, at the end of the play reveals his true identity and each find their own individual destinies with their sweethearts, Cecily and Gwendolin. We leave it as an exercise for the reader to identify these four characters with the subject matter of this paper...in any event, all's well that ends well!

As in years past, I would like once again to thank friends and colleagues at the Instituto de Investigaciones en Materiales for their hospitality, particularly at the present time for hosting his sabbatical. Special thanks also go to the staff of the IBM Almaden Research Center library for consultation on the theatrical literature.

## REFERENCES

[1] F. Marsiglio, 12th Winter Meeting on Low Temperature Physics, Vista Hermosa, Morelos, Mexico, Progress in High Temperature Superconductivity (World Scientific, Singapore, 1991), present volume.

[2] J. E. Hirsch, 10th Winter Meeting on Low Temperature Physics,

Cocoyoc, Morelos, Mexico, Progress in High Temperature Superconductivity (World Scientific, Singapore, 1989), p. 43; J. E. Hirsch, in "Electronic Structure and Mechanisms for High Temperature Superconductivity" (Proceedings of the Miami Workshop, January 3-9, 1991), ed. J. Ashkenazi and G. Vezzoli, Plenum, New York, 1991.

[3] I. S. Shaplygin, B. G. Kakhan and V. B. Lazerev, *Rus. J. of Inorg. Chem.* 24, 820 (1979).

[4] Y. Tokura, H. Takagi and S. Uchida, *Nature* 337, 345 (1989).

[5] H. Takagi, S. Uchida and Y. Tokura, *Phys. Rev. Letters* 62, 1197 (1989).

[6] M. E. Lopez-Morales, R. J. Savoy and P. M. Grant, *Solid State Commun.* 71, 1079 (1989).

[7] Z. Z. Wang, D. A. Brawner, T. R. Chein, N. P. Ong, J. M. Tarascon and E. Wang, *Physica C* 162-164, 1013 (1989); Recent measurements of the Hall effect on superconducting n-type single crystals consistently yield negative Hall coefficients down to  $T_c$  (R. L. Greene, private communication).

[8] M. E. Lopez-Morales, R. J. Savoy and P. M. Grant, *J. Mater. Res.* 5, 2041 (1990).

[9] J. M. Tranquada, S. M. Heald, A. R. Moodenbaugh, G. Liang and M. Croft, *Nature* 337, 720 (1989); A. Grassmann, J. Schloetter, J. Stroebel, M. Klauda, R. L. Johnson and G. Saemann-Ischenko, *Physica C* 162-164, 1383 (1989).

[10] P. M. Grant, *Adv. Mater.* 2, 232 (1990).

- [11] T. Siegrist, S. M. Zahurak and D. W. Murphy, *Nature* 334, 171 (1988).
- [12] M. G. Smith, A. Manthiram, J. Zhou, J. B. Goodenough and J. T. Markert, *Nature* 351, 549 (1991).
- [13] R. J. Cava, B. Batlogg, R. B. van Dover, J. J. Kajewski, J. V. Waszczak, R. M. Fleming, W. F. Peck, Jr., L. W. Rupp, Jr., P. Marsh, A. C. W. P. James and L. F. Scheemeyer, *Nature* 345, 602 (1990).
- [14] R. J. Cava, T. Siegrist, B. Hesse, J. J. Krajewski, W. F. Peck, Jr., B. Batlogg, H. Takagi, J. V. Waszczak, L. F. Scheemeyer and H. W. Zandbergen, *Physica C*, in press; R. J. Cava,  $M_2S$  - HTSC III, Kanazawa, Japan, 22-26 July 1991.

Following are page proofs.

## THE IMPORTANCE OF BEING N-TYPE

P. M. GRANT\*

*Instituto de Investigaciones en Materiales, Universidad Nacional Autónoma de México, Apartado Postal 70-360, 04510 México, D.F.*

### ABSTRACT

Perhaps the most significant recent advance in high- $T_c$  materials has been the discovery of superconductivity in the  $\text{Nd}_2\text{CuO}_4$ -based systems. Certain solid state measurements, primarily Hall effect and thermoelectric power, would suggest transport in the normal state, and by implication also the superconducting state, of these new compounds is dominated by electrons rather than holes, the latter having been the case for all previous high- $T_c$  materials. This discovery of charge reciprocity is potentially important for guiding theoretical developments. However, the n-type compounds differ significantly in both their preparation and structural properties from their p-type antecedents, and display important non-reciprocal behavior in doping and oxidation-reduction treatments. In this article, we will review the present situation regarding n-type high temperature superconductors, paying particular attention to the aforementioned differences from the hole carrier materials and the impact such differences may have on extending the number of presently known n-type systems.

\* *On sabbatical from the IBM Almaden Research Center, San Jose, California, USA.*

Electron-hole duality has long been a mainstay in the modern theory of the electronic properties of semiconductors and metals. This was not always the case. In the 19th century, it was very perplexing to discover that some metals had a positive Hall coefficient, and even worse, many yielded Seebeck coefficients opposite in sign to that of the Hall effect. These questions were resolved in large part with the arrival of the quantum theory of solids in the 1930s and 40s, although some, especially those involving the Seebeck effect in noble metals, still seek complete answers today.

As regards superconductivity, conventional BCS theory makes no distinction between holes and electrons. Indeed, the number of superconductors with a hole-like normal state is manifold, and there is some feeling (see the paper by Marsiglio in this conference [1], and the one presented by Hirsch at a preceding meeting in this series held in Cocoyoc [2]) that this may be a universal property of all superconductors, requiring, perhaps, extensive departures from the BCS concept.

All early high temperature superconductors displayed the usual aspects of possessing positive charge normal state carriers...positive Hall and Seebeck response, reinforced by simply adding up the excess charge expected from conventional ionic charge states on the constituent ions. From the very beginning attempts were made (all unsuccessful, and none, to the knowledge of the author, ever published) to discover electron analogs to the  $\text{La}_2\text{CuO}_4$  system. Those participated in by the author included tetravalent ion substitutions (e.g., Pb and Sn) and the use of various vanadate hosts (the electron equivalent of  $\text{Cu}^{++}$ ). In fact, the literature already included several strong hints as to how to accomplish this objective. A paper by Shaplygin, et al. [3], preceding by

several years the high-Tc discoveries, contained a table of Seebeck coefficients for the isoelectronic series  $\text{Ln}_2\text{CuO}_4$ , where Ln was La or a trivalent rare earth ion. All, except  $\text{La}_2\text{CuO}_4$ , were negative in sign at room temperature. It was thus natural to attempt substituting a tetravalent cation on the lanthanide position in one of these latter compounds. This step was taken by Tokura and co-workers [4] in late 1988 when they synthesized Ce-doped  $\text{Nd}_2\text{CuO}_4$ . However, metallic conductivity, and hence superconductivity, was not obtained until a further crucial step was taken...it was necessary to anneal the samples in a reducing atmosphere and then quench to room temperature. It was also found that Th-doping worked and that  $\text{Sm}_2\text{CuO}_4$  and  $\text{Pr}_2\text{CuO}_4$  also provided a suitable host system. Charge counting assuming normal ionic state values resulted in net negative charge per unit cell, which was upheld by temperature dependent measurements of the Hall constant [5] and thermoelectric power [6], both revealing a negative sign down to the onset of superconductivity. Other early experiments [7] indicated that the signs of both these effects went positive before the superconducting state was reached, thus suggesting that these new materials were hole superconductors as well, but it is now nearly certain that this behavior was caused by inhomogeneous doping due to the slow diffusion of the tetravalent species in the host material in the case of ceramics [8], or, non-uniform distribution of oxygen vacancies in the case of single crystals.

The consistently negative sign of the thermopower and Hall constant, supported by the expected excess negative charge from the ionic states involved, are compelling, although not conclusive, arguments for electron pairing in the superconducting state. There is always the possibility, albeit small, that multiband effects and pathological Fermi surface topology could complicate this straightforward interpretation. Both of these carrier-sign-sensitive measurements in metals involve complex

integrations of dyad products of the group velocity (i.e., dispersion in the energy vs. quasimomentum band structure) over the Fermi surface which in principle could be dominated by small regions of high curvature. This situation is, unlikely given the small carrier concentration in high-T<sub>c</sub> superconductors relative to most other metals, so that the carriers are very likely confined to band extrema where the effective mass approximation holds and an unambiguous sign can be attached to the carriers. More uncertain is whether band overlap occurs, resulting in electron-hole pockets in which electrons dominate the normal state, yet holes still form the paired state. If, in the p-type compounds, holes reside in O 2p - Cu 3d hybridized valence bands, and, for the n-type case the electrons transport in the Cu 4s conduction band, the most consistent interpretation of core level spectra obtained for each material class [9], then band overlap is improbable given the usual magnitude of the energy gap in transition metal oxides between these two bands. Thus it is prudent to conclude at the present time that charge reciprocity in high temperature superconductors exists and must be accommodated by any theory put forward to describe these materials.

The existence of charge symmetry does not imply structural or chemical symmetry between n- and p-type high-T<sub>c</sub> compounds such as found, for example, in covalent semiconductors. Important differences in both crystallographic structure and chemical preparation arise. In Fig. 1 we compare unit cells of the two compounds and the respective doping sites. The principal difference is the coordination of the lanthanide cation with respect to the nearest neighbor oxygen anions. In La<sub>2</sub>CuO<sub>4</sub>, there is Cu-O plane apical oxygen which in Nd<sub>2</sub>CuO<sub>4</sub> moves to its observed position on the unit cell face from its former edge location. Apparently this lack of an apical oxygen is the key structural feature which permits electron delocalization. The host structure choice



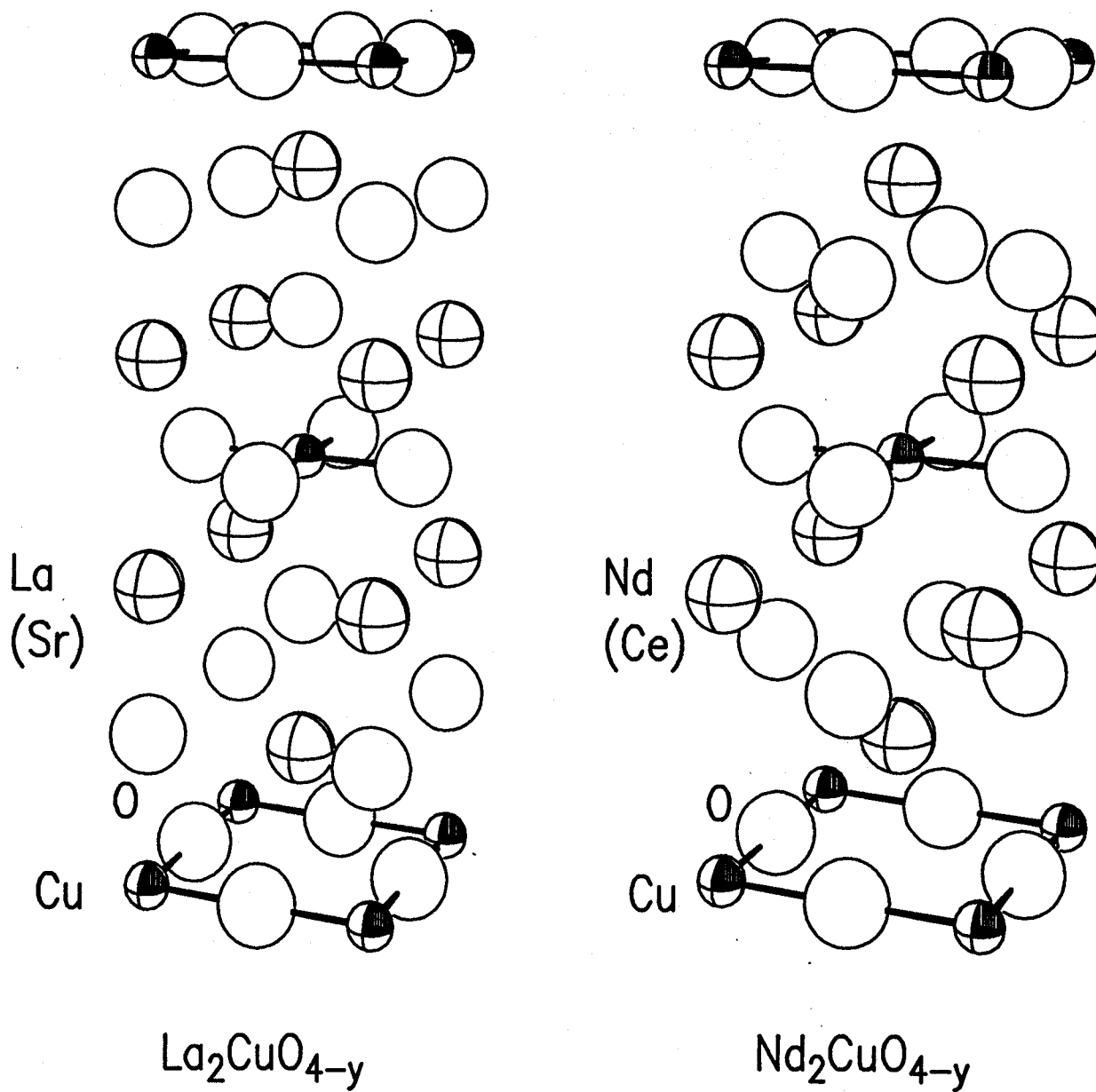


Fig. 1. Unit cells of  $\text{La}_2\text{CuO}_{4-y}$  and  $\text{Nd}_2\text{CuO}_{4-y}$ .  
Note the absence of apical oxygens  
in the latter structure.

possessing this feature for potential n-type superconductivity has been largely restricted to the  $\text{Nd}_2\text{CuO}_4$  class, inasmuch as all other Cu-O planar perovskite structures have at least one apically coordinated plane and appear p-type, with any accompanying non-apical planes "dead" with regard to transport; e. g., the middle Cu-O layer in the 2223 structure. As has been pointed out [10], the single exception was the "infinite layer" compound  $\text{Ca}_{0.86}\text{Sr}_{0.14}\text{CuO}_2$ , which was discovered by Siegrist [11] and coworkers in mid-1988. Recently, two structural isomorphs,  $\text{Sr}_{0.86}(\text{Nd,Pr})_{0.14}\text{CuO}_2$ , were synthesized by Smith, et al., [12] and both were found to have superconducting transitions near 40K. By unit cell charge count, the substitution of trivalent Nd or Pr should result in excess negative charge and thus an n-type normal state. However, it was found by the authors of ref. [12] that their preliminary measurements of the Hall constant were positive, a finding they speculate is due to minority phases in their samples. If subsequent research reveals  $R_H < 0$ , especially near the transition temperature, as additional confirmation of charge counting, then the discovery of n-type superconductivity at 40K is an important advance in at least two aspects. Firstly, unlike the  $\text{Nd}_2\text{CuO}_4$ -based systems, the "infinite layer" compounds do not require direct reduction processing. On the contrary, both the precursor compounds and the final phase are obtained via high pressure oxygen processing. However, when more is known about the microstructure, it may turn out that oxygen vacancies indeed exist. Secondly, and most critically, the discovery of this new high-Tc material demonstrates that transition temperatures as high as those obtained in Sr-doped  $\text{La}_2\text{CuO}_4$  can be reached in systems without apical oxygens, thus proving that the existence of this ion is not a necessary condition for truly high temperature superconductivity.

Where do we go from here? Superconductivity at 60K has been reported in a "double layer" variant of doped- $\text{La}_2\text{CuO}_4$  [13]. In

fact, a whole series of multiple layer variants seem possible [14]. A suggestive path would be to attempt the synthesis of appropriately n-doped T' isomorphs where trivalent rare earth replacement of the lanthanum would move the apical oxygen to the cell edges as in  $\text{Nd}_2\text{CuO}_4$ . Only time and clever chemistry will tell.

The astute reader may recognize the title of this article as a pun, albeit far-fetched, on Oscar Wilde's famous 1895 play, "The Importance of Being Earnest." Elgernon, who originally pretended to be Jack, at the end of the play reveals his true identity and each find their own individual destinies with their sweethearts, Cecily and Gwendolin. We leave it as an exercise for the reader to identify these four characters with the subject matter of this paper...in any event, all's well that ends well!

As in years past, I would like once again to thank friends and colleagues at the Instituto de Investigaciones en Materiales for their hospitality, particularly at the present time for hosting his sabbatical. Special thanks also go to the staff of the IBM Almaden Research Center library for consultation on the theatrical literature.

#### REFERENCES

[1] F. Marsiglio, 12th Winter Meeting on Low Temperature Physics, Vista Hermosa, Morelos, Mexico, Progress in High Temperature Superconductivity (World Scientific, Singapore, 1991), present volume.

[2] J. E. Hirsch, 10th Winter Meeting on Low Temperature Physics,

Cocoyoc, Morelos, Mexico, Progress in High Temperature Superconductivity (World Scientific, Singapore, 1989), p. 43; J. E. Hirsch, in "Electronic Structure and Mechanisms for High Temperature Superconductivity" (Proceedings of the Miami Workshop, January 3-9, 1991), ed. J. Ashkenazi and G. Vezzoli, Plenum, New York, 1991.

[3] I. S. Shaplygin, B. G. Kakhan and V. B. Lazerev, Rus. J. of Inorg. Chem. 24, 820 (1979).

[4] Y. Tokura, H. Takagi and S. Uchida, Nature 337, 345 (1989).

[5] H. Takagi, S. Uchida and Y. Tokura, Phys. Rev. Letters 62, 1197 (1989).

[6] M. E. Lopez-Morales, R. J. Savoy and P. M. Grant, Solid State Commun. 71, 1079 (1989).

[7] Z. Z. Wang, D. A. Brawner, T. R. Chein, N. P. Ong, J. M. Tarascon and E. Wang, Physica C 162-164, 1013 (1989); Recent measurements of the Hall effect on superconducting n-type single crystals consistently yield negative Hall coefficients down to  $T_c$  (R. L. Greene, private communication).

[8] M. E. Lopez-Morales, R. J. Savoy and P. M. Grant, J. Mater. Res. 5, 2041 (1990).

[9] J. M. Tranquada, S. M. Heald, A. R. Moodenbaugh, G. Liang and M. Croft, Nature 337, 720 (1989); A. Grassmann, J. Schloetter, J. Stroebel, M. Klauda, R. L. Johnson and G. Saemann-Ischenko, Physica C 162-164, 1383 (1989).

[10] P. M. Grant, Adv. Mater. 2, 232 (1990).

[11] T. Siegrist, S. M. Zahurak and D. W. Murphy, Nature 334, 171 (1988).

[12] M. G. Smith, A. Manthiram, J. Zhou, J. B. Goodenough and J. T. Markert, Nature 351, 549 (1991).

[13] R. J. Cava, B. Batlogg, R. B. van Dover, J. J. Krajewski, J. V. Waszczak, R. M. Fleming, W. F. Peck, Jr., L. W. Rupp, Jr., P. Marsh, A. C. W. P. James and L. F. Schneemeyer, Nature 345, 602 (1990).

[14] R. J. Cava, T. Siegrist, B. Hessen, J. J. Krajewski, W. F. Peck, Jr., B. Batlogg, H. Takagi, J. V. Waszczak, L. F. Scheemeyer and H. W. Zandbergen, Physica C, in press; R. J. Cava,  $M_2S$  - HTSC III, Kanazawa, Japan, 22-26 July 1991.